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# Winds of hot subdwarfs: metallicity is the key

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# Stellar wind of hot stars

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Puls et al. (2008):

- supersonic flow from hot stars
- accelerated due to the absorption of radiation mainly in the resonance lines of heavier elements (C, N, O, or Fe)
- found in OBA stars, WR stars, CSPNs, and sdO stars

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- wind mass-loss rate  $\dot{M}$  depends on
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- the terminal velocity  $v_\infty$  proportional to the escape velocity  $v_{\text{esc}}$

$$v_\infty \sim v_{\text{esc}}$$

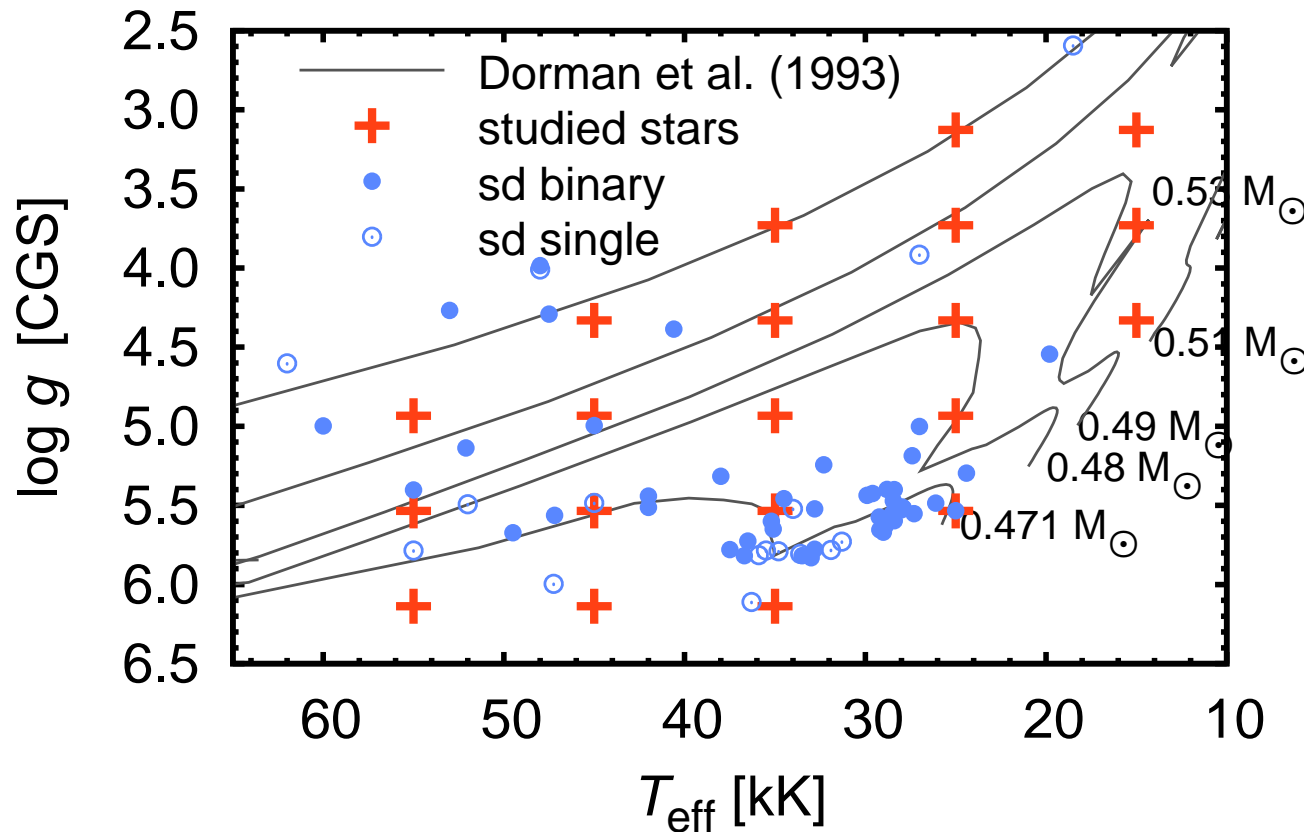
# NLTE wind models

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(Krtička & Kubát 2010)

- spherically symmetric stationary wind models
  - level populations derived from kinetic equilibrium (NLTE) equations
  - radiative force calculated using the comoving frame (CMF) radiative transfer equation
  - wind density, velocity and temperature structure derived from hydrodynamical equations
- ⇒ enable to predict wind mass-loss rate  $\dot{M}$ , and terminal velocity  $v_\infty$

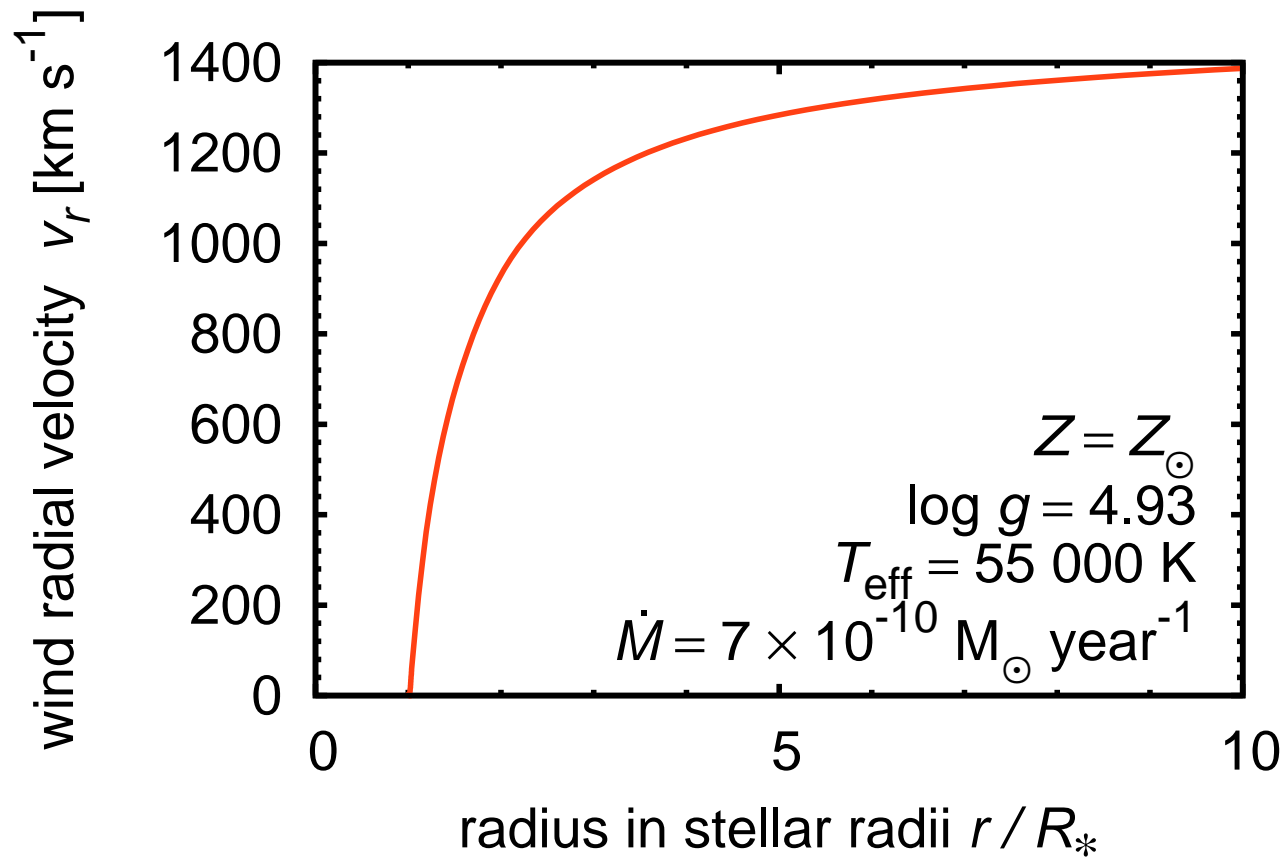
# Model stars: $\log g$ vs. $T_{\text{eff}}$



- stellar mass  $M = 0.5 M_{\odot}$
- metallicity  $Z$ :  $0.1 Z_{\odot}$ ,  $Z_{\odot}$ ,  $10 Z_{\odot}$
- helium content  $N(\text{He})/N(\text{H})$ : 0.085

# Example of wind model

- wind radial velocity

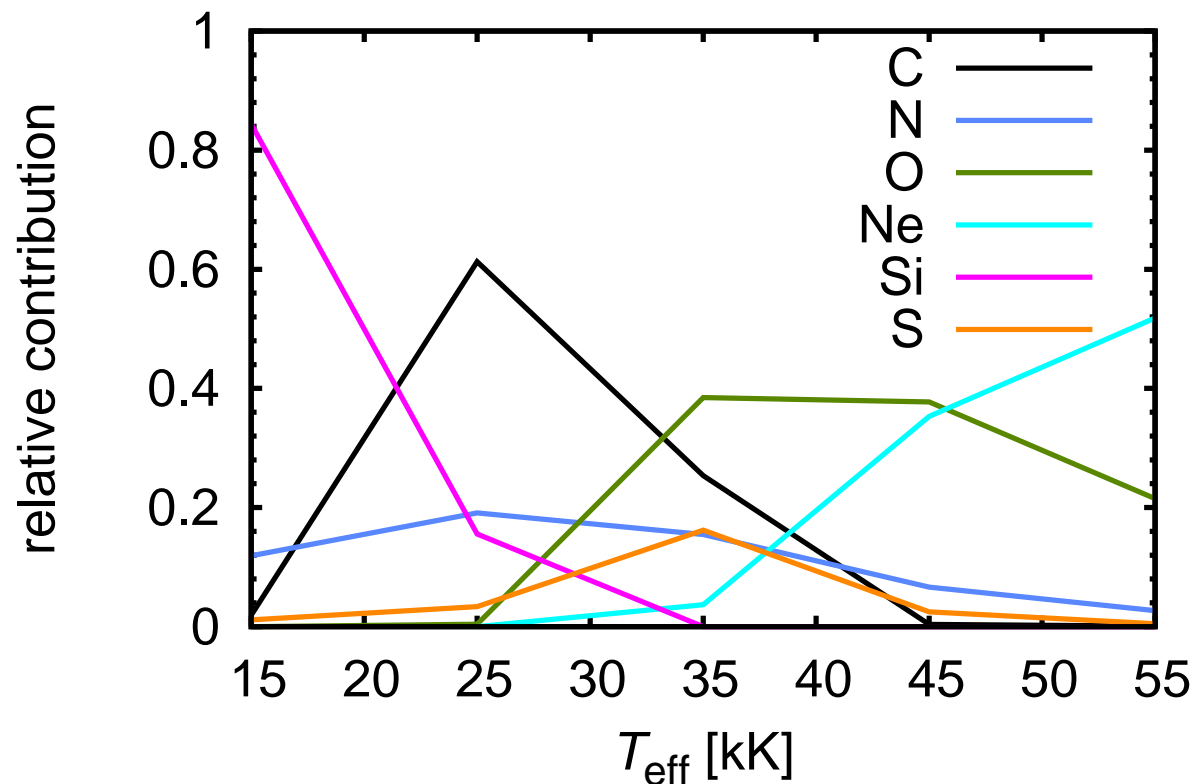


- low luminosity  $\Rightarrow$  low mass-loss rate



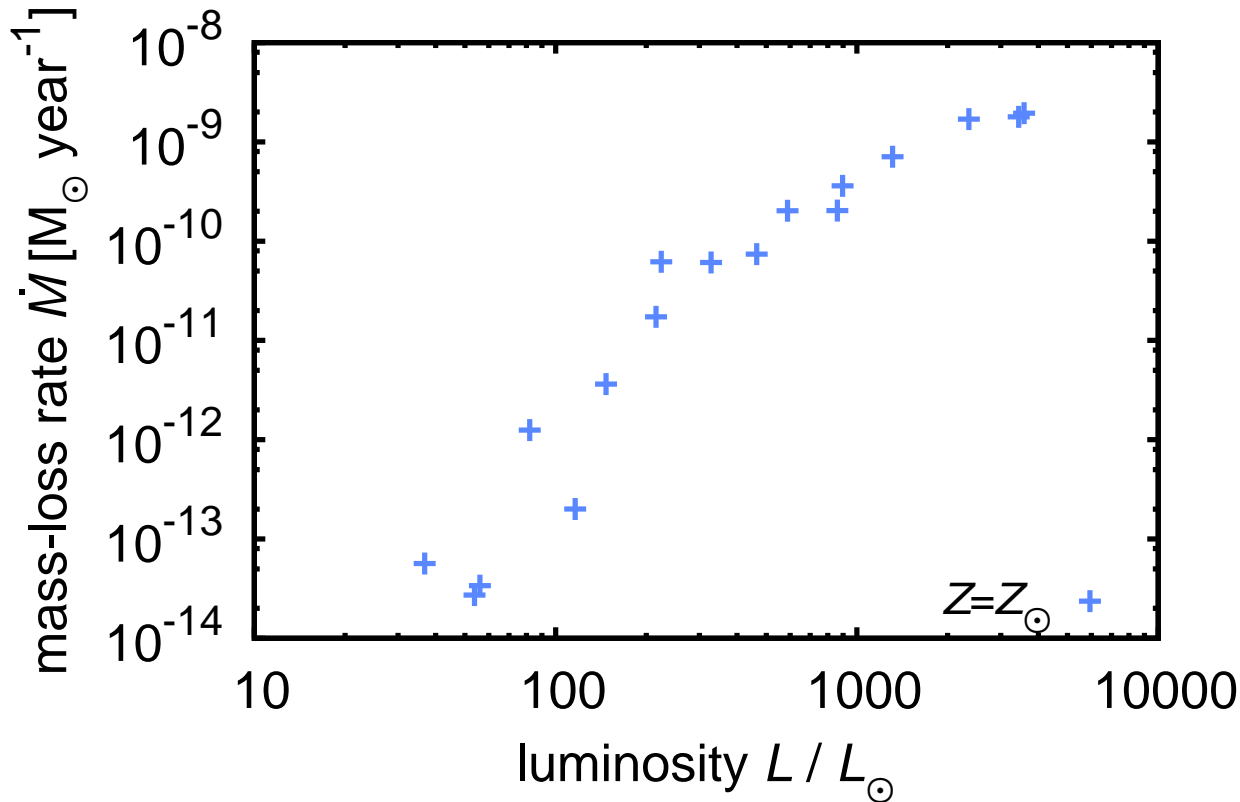
# Contribution of different elements

- relative contribution of individual elements to the radiative force for  $Z = Z_{\odot}$



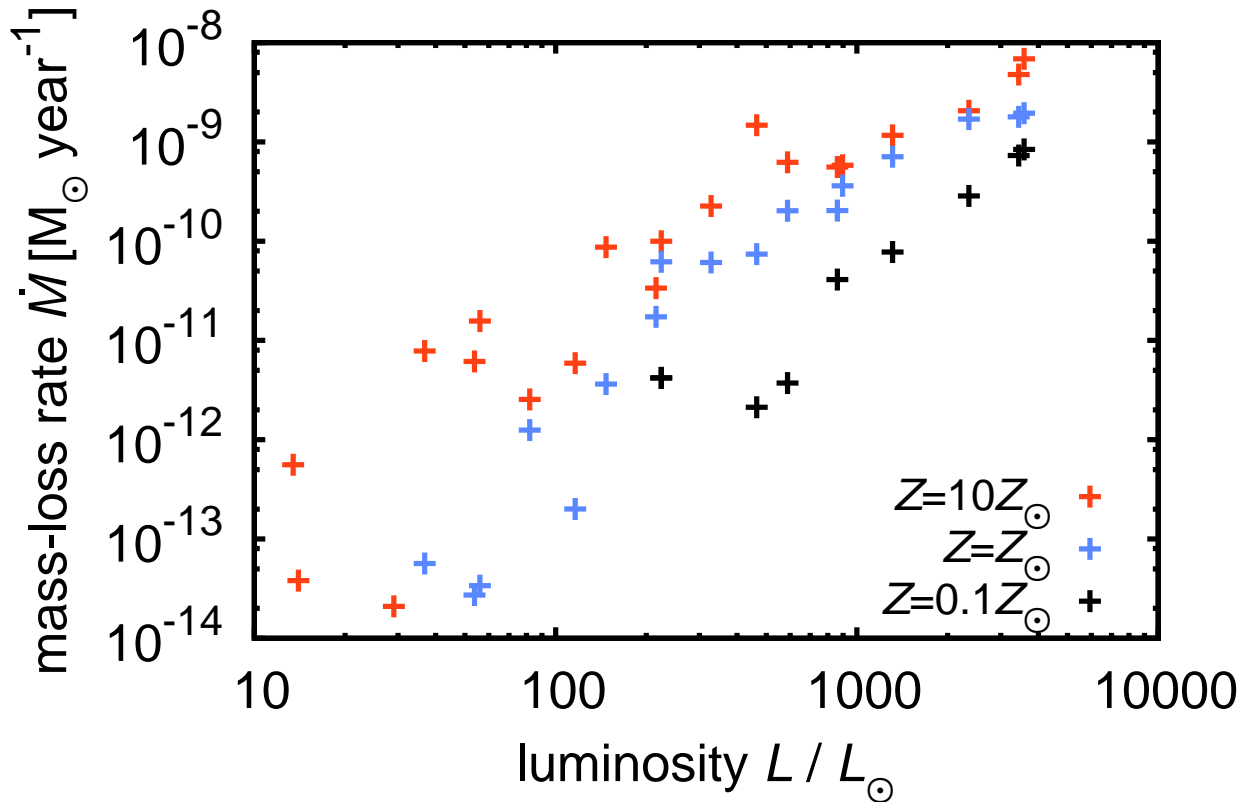
- different elements contribute at different  $T_{\text{eff}}$
- mass-loss rate sensitive to the composition

# Wind mass-loss rate



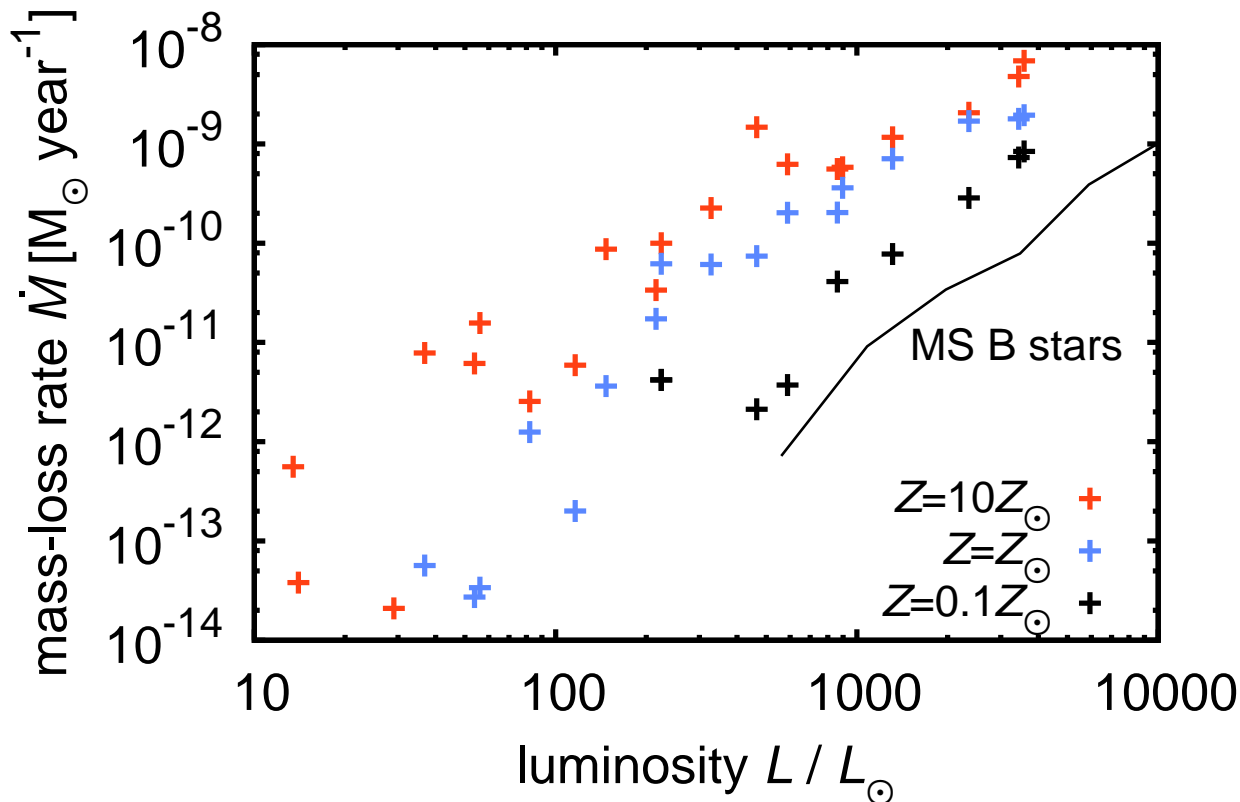
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# Wind mass-loss rate



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# Wind mass-loss rate



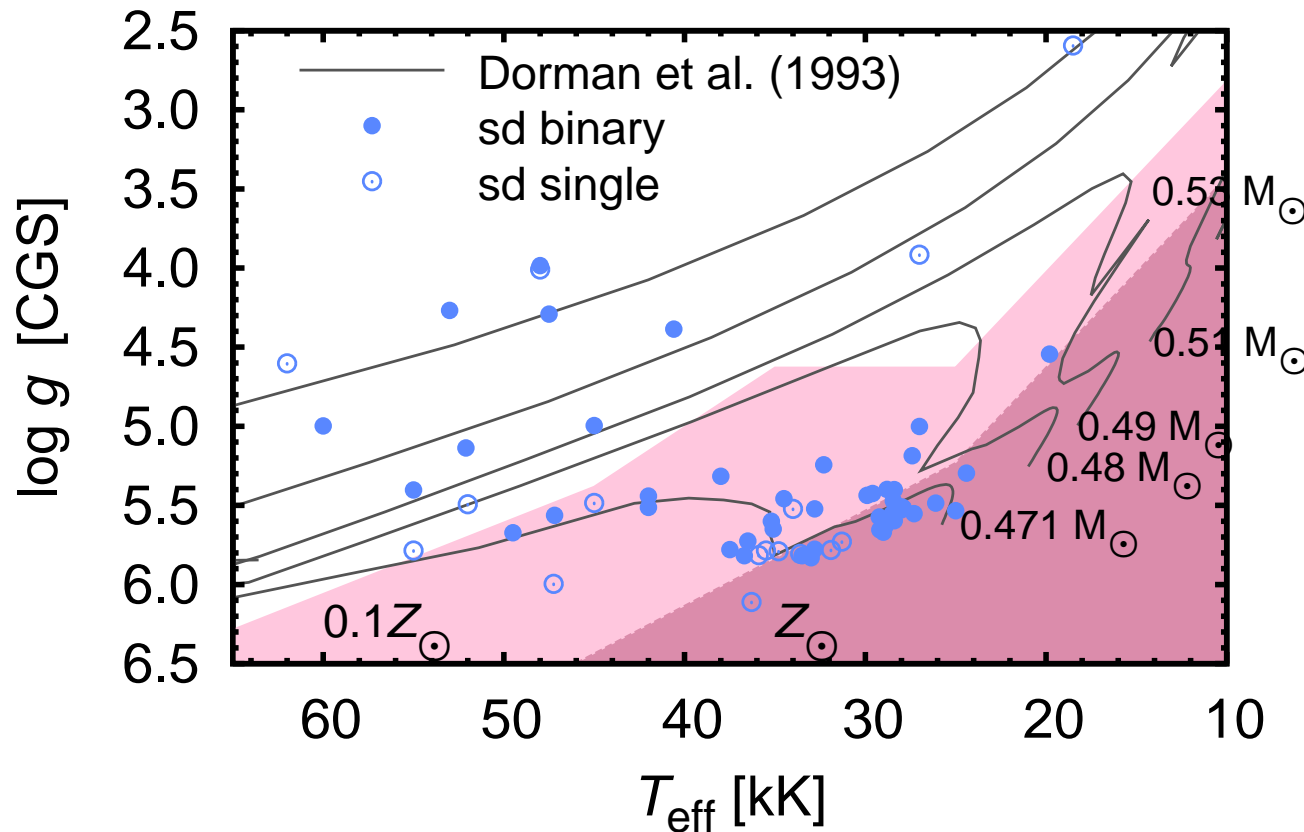
- mass-loss rate scales with stellar luminosity
- mass-loss rate depends on metallicity
- higher than for main-sequence B stars with the same luminosity

# Wind mass-loss rate

- mass-loss rate recipe

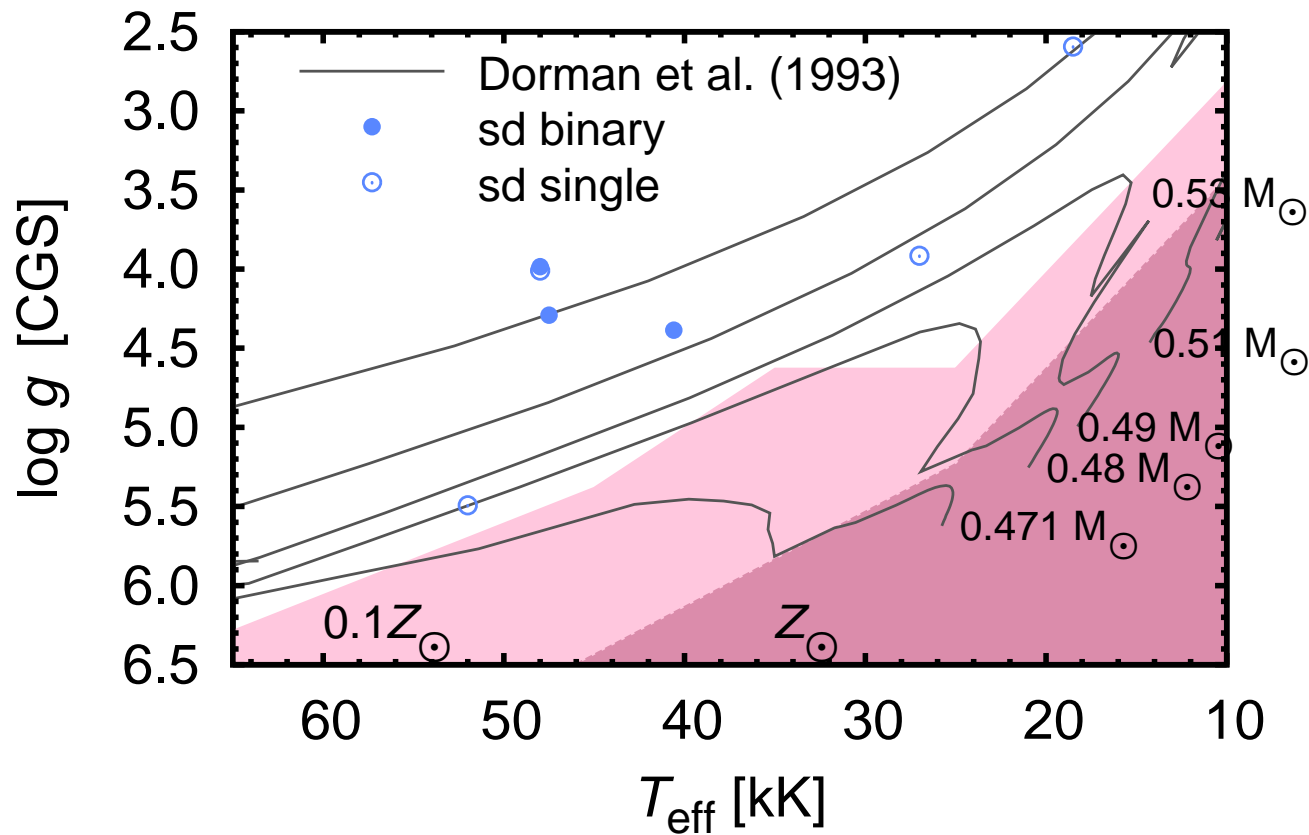
$$\begin{aligned} \log\left(\frac{\dot{M}}{1 M_{\odot} \text{ year}^{-1}}\right) &= -12.6 + 1.4 \log \frac{Z}{Z_{\odot}} + \\ &+ \left(3.8 - 1.4 \log \frac{Z}{Z_{\odot}}\right) \log\left(\frac{L}{10^2 L_{\odot}}\right) + \\ &+ \left(-1.1 + 0.6 \log \frac{Z}{Z_{\odot}}\right) \log^2\left(\frac{L}{10^2 L_{\odot}}\right) + \\ &+ 1.1 \log\left(\frac{T_{\text{eff}}}{10^4 \text{ K}}\right) \end{aligned}$$

# Wind domains: $\log g$ vs. $T_{\text{eff}}$



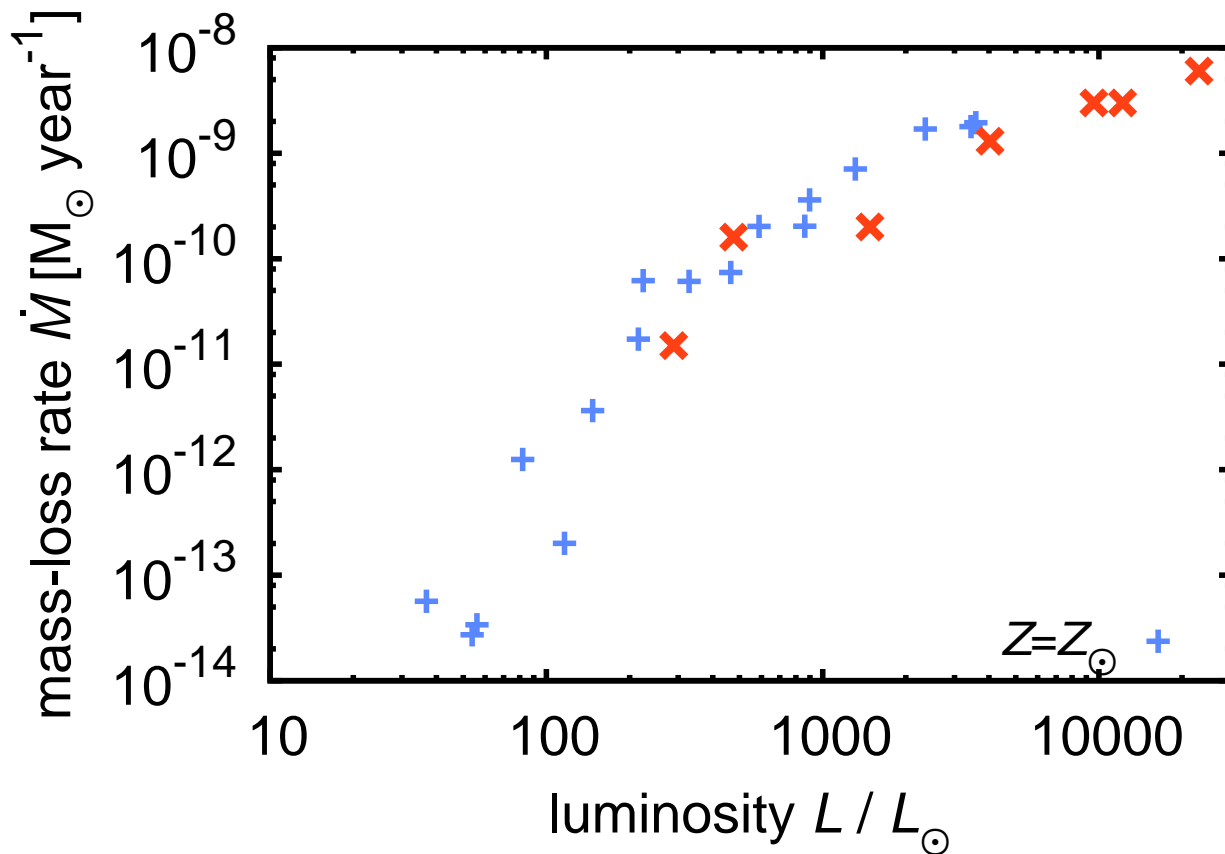
- wind possible only in white areas ( $g^{\text{rad}} > g$ )
- forbidden (red) region with  $g^{\text{rad}} < g$

# Stars with winds: $T_{\text{eff}}$ vs. $\log g$



- stars with observed wind in the allowed region

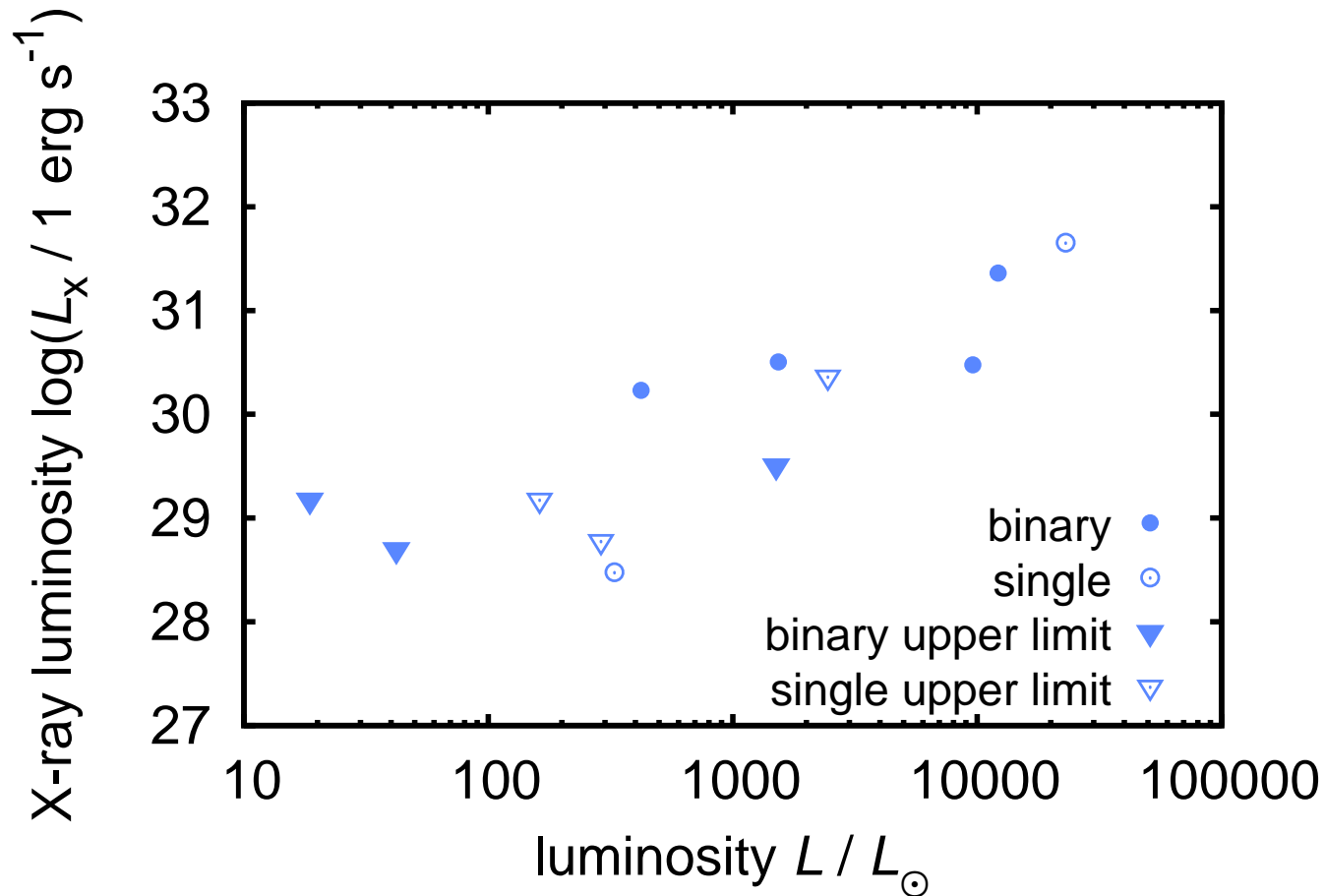
# Test against observations: $\dot{M}$



- mass-loss rates: Jeffery & Hamann (2010), Lanz et al. (1997), Gruschinske et al. (1983)

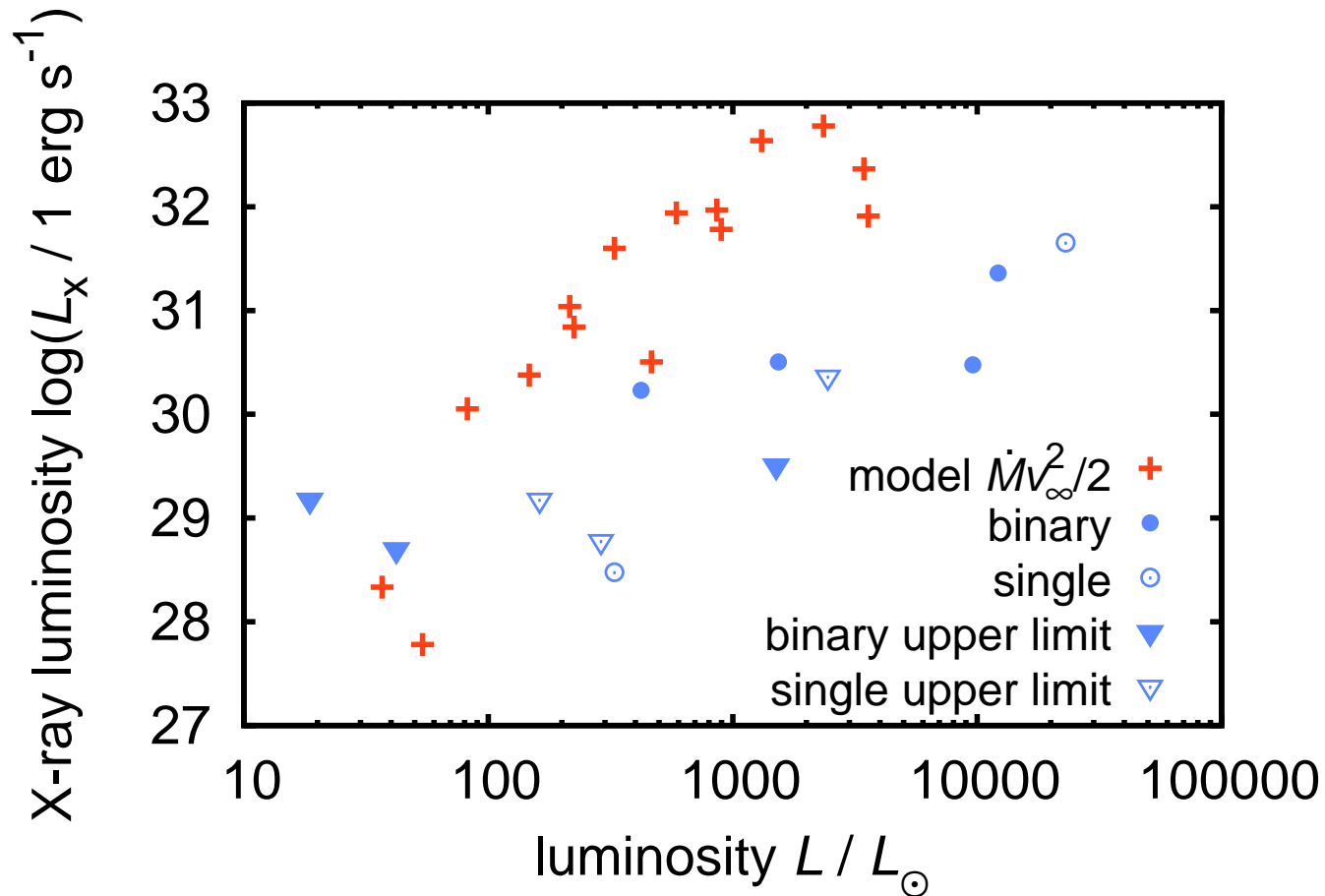


# Test against observations: $L_x$



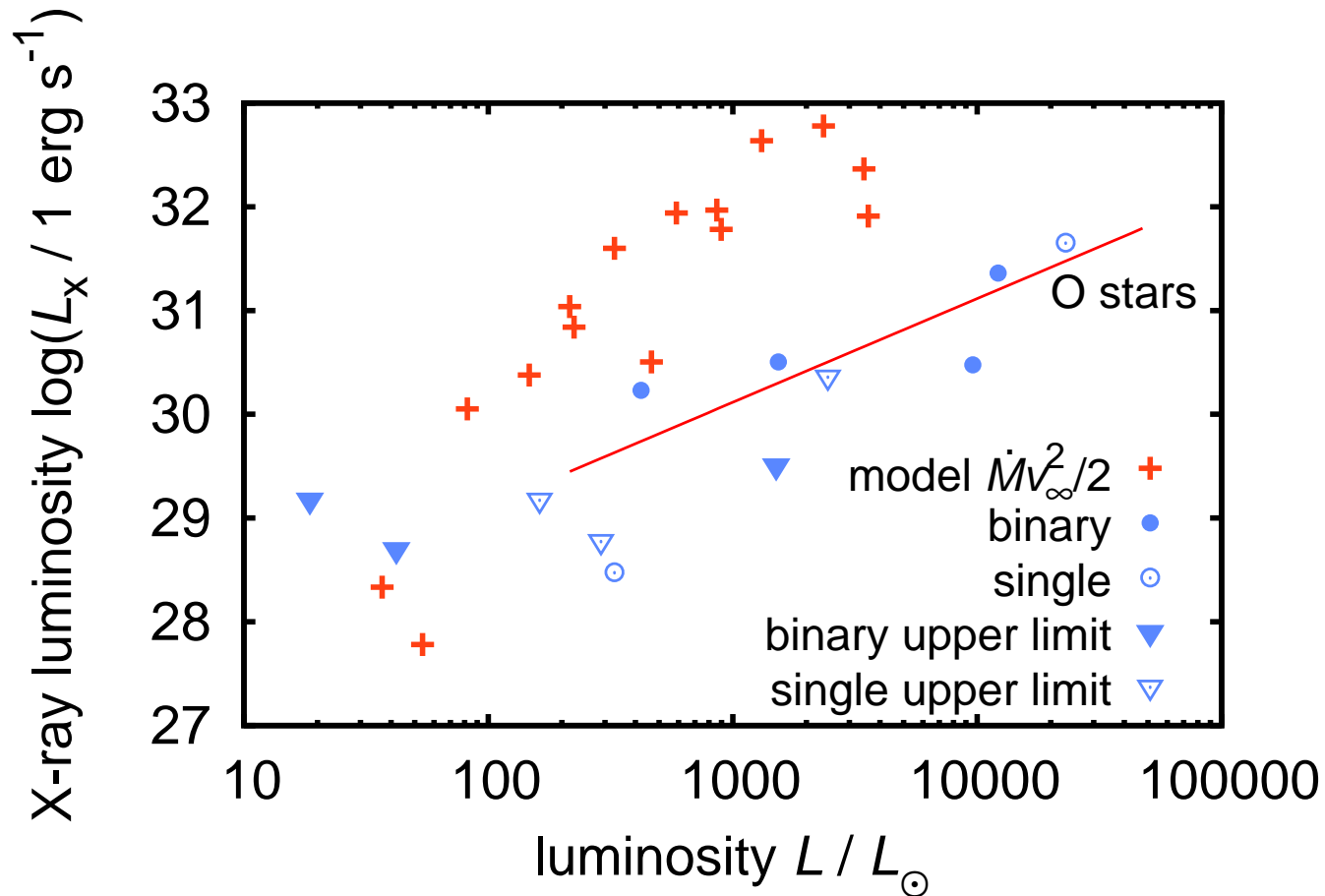
- X-ray luminosities: La Palombara et al. (2012), Montez et al. (2010), Mereghetti et al. (2013), La Palombara et al. (2014)

# Test against observations: $L_x$



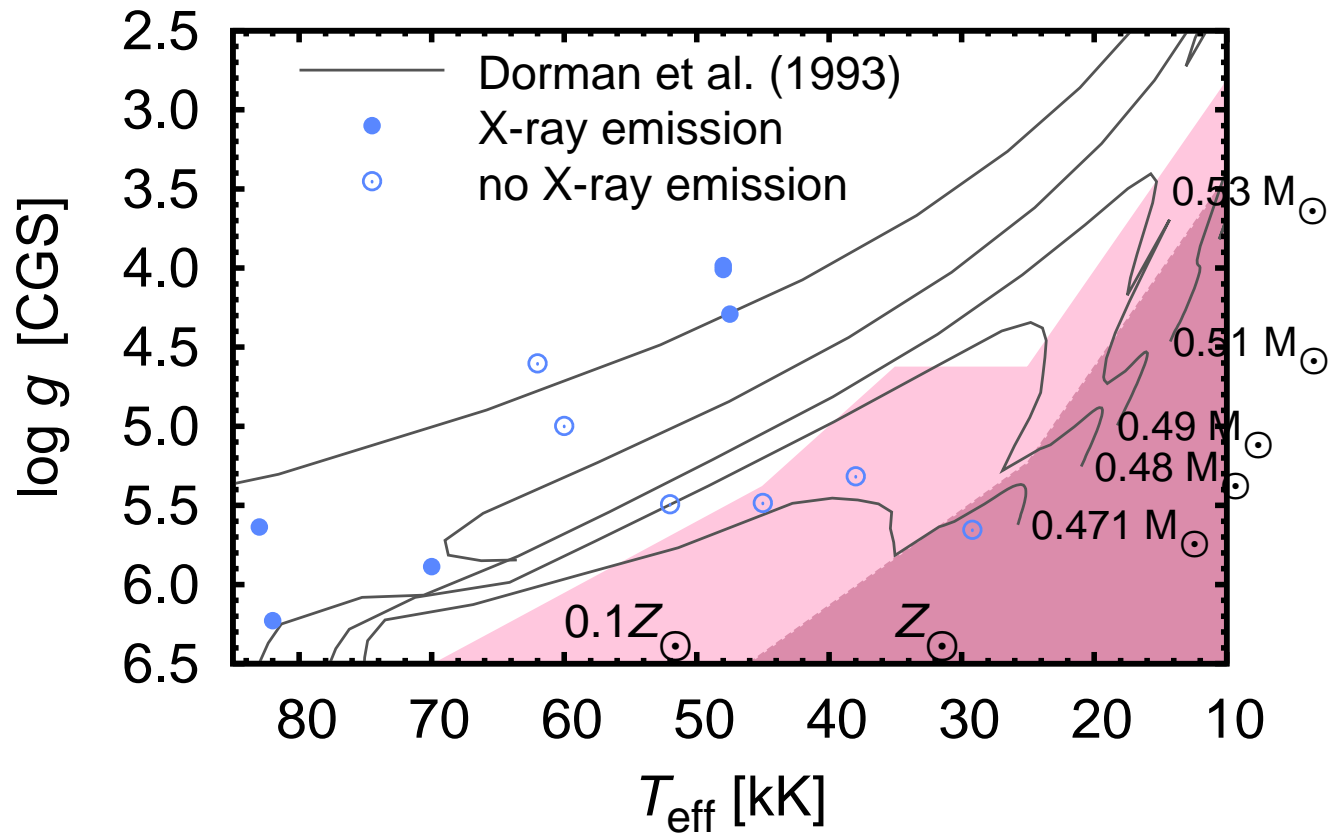
- X-rays may be produced in the wind  
( $\frac{1}{2} \dot{M} v_\infty^2 > L_x$ )

# Test against observations: $L_x$



- X-ray luminosity corresponds to O stars (Nazé 2009)

# X-ray emission: $\log g$ vs. $T_{\text{eff}}$



- X-ray emission only in the wind domain

# No X-ray emission: CD $-30^\circ$ 11223

- CD  $-30^\circ$  11223: sdB+WD binary
- Bondi-Hoyle-Lyttleton accretion of the wind on WD companion
- X-ray luminosity ( $a$  is the orbital separation)

$$L_x = G \frac{M_{\text{WD}} \dot{M}_{\text{acc}}}{R_{\text{WD}}} = \frac{1}{4} \frac{G M_{\text{WD}}}{R_{\text{WD}}} \dot{M} \left( \frac{r_{\text{acc}}}{a} \right)^2,$$

the accretion radius is  $r_{\text{acc}} = 2GM_{\text{WD}}/v^2$

(Mereghetti et al. 2014)

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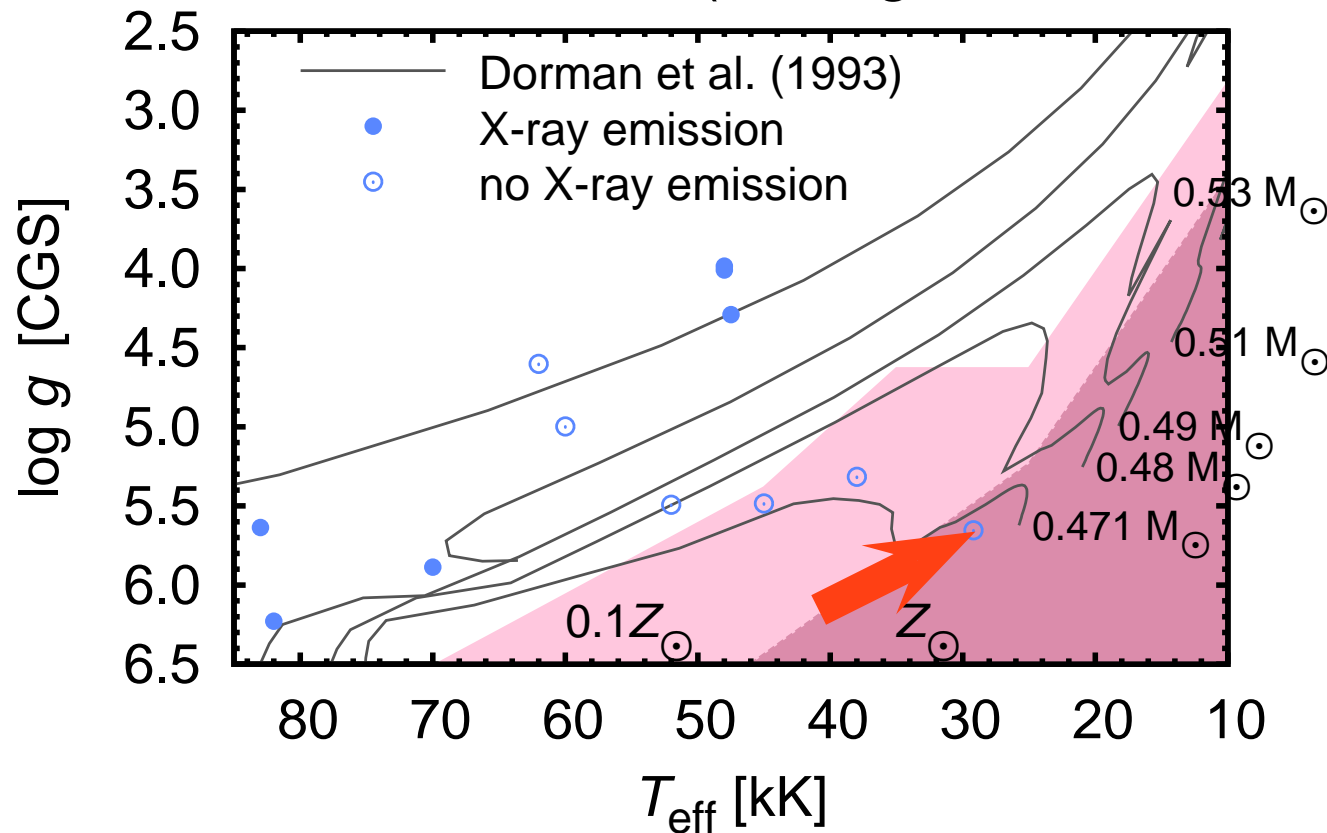
- no X-rays detected from CD  $-30^\circ$  11223:  
 $L_x < 1.5 \times 10^{29} \text{ erg s}^{-1}$
- upper limit for the mass-loss rate  
 $\dot{M} < 3 \times 10^{-13} M_\odot \text{ year}^{-1}$

(Mereghetti et al. 2014)

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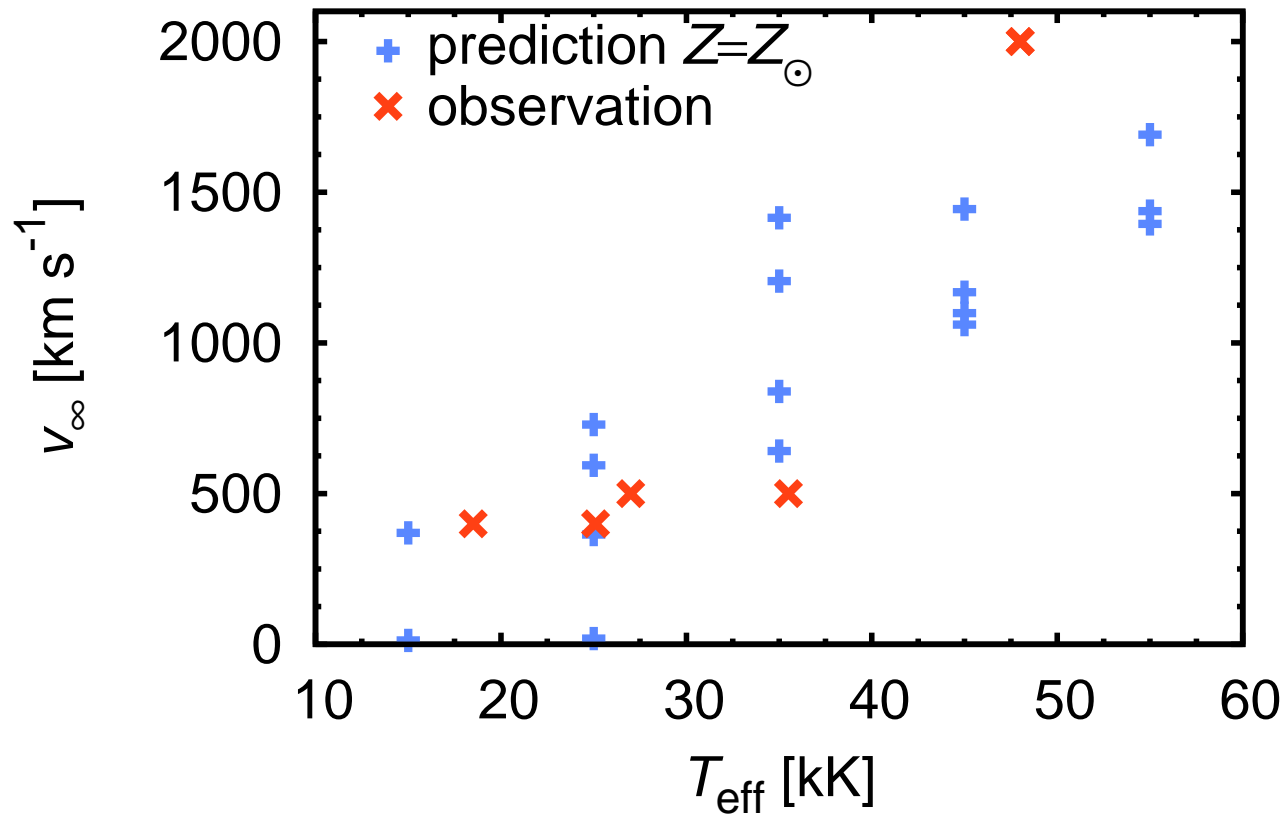
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- $\dot{M} < 3 \times 10^{-13} M_\odot \text{ year}^{-1}$

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- no wind from the simulations

# Test against observations: $v_\infty$



- wind terminal velocity  $v_\infty$  proportional to  $v_{\text{esc}}$
- decreases with  $T_{\text{eff}}$  due to decrease of  $v_{\text{esc}}$
- observations: Jeffery & Hamann (2010)



# Cherry on top: $\varphi$ Per

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- member of a rare group of sdO+Be binaries (Gies et al. 1998, Koubský et al. 2014)
- sdO star orbits in the disk of Be star



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- member of a rare group of sdO+Be binaries (Gies et al. 1998, Koubský et al. 2014)
- sdO star orbits in the disk of Be star
- sdO star mass-loss rate  $1.1 \times 10^{-9} M_{\odot} \text{ year}^{-1}$
- Be star disk mass-loss rate  $10^{-9} M_{\odot} \text{ year}^{-1}$  (Granada et al. 2013)
- the disk and wind collide: momentum condition

$$\rho_{\text{disk}} v_{\text{disk}}^2 = \rho_{\text{wind}} v_{\text{wind}}^2$$

- this occurs at  $0.1 - 0.3 a$  (orbital separation)
- ⇒ radiative and mechanical interaction

# Conclusions: subdwarf winds

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- driven by lighter elements (similarly to early B stars)
- stars with low effective temperatures and high surface gravities do not have any wind
- metallicity is one of the key wind parameters
- fair agreement with observations (mass-loss rates, X-ray luminosities, terminal velocities)