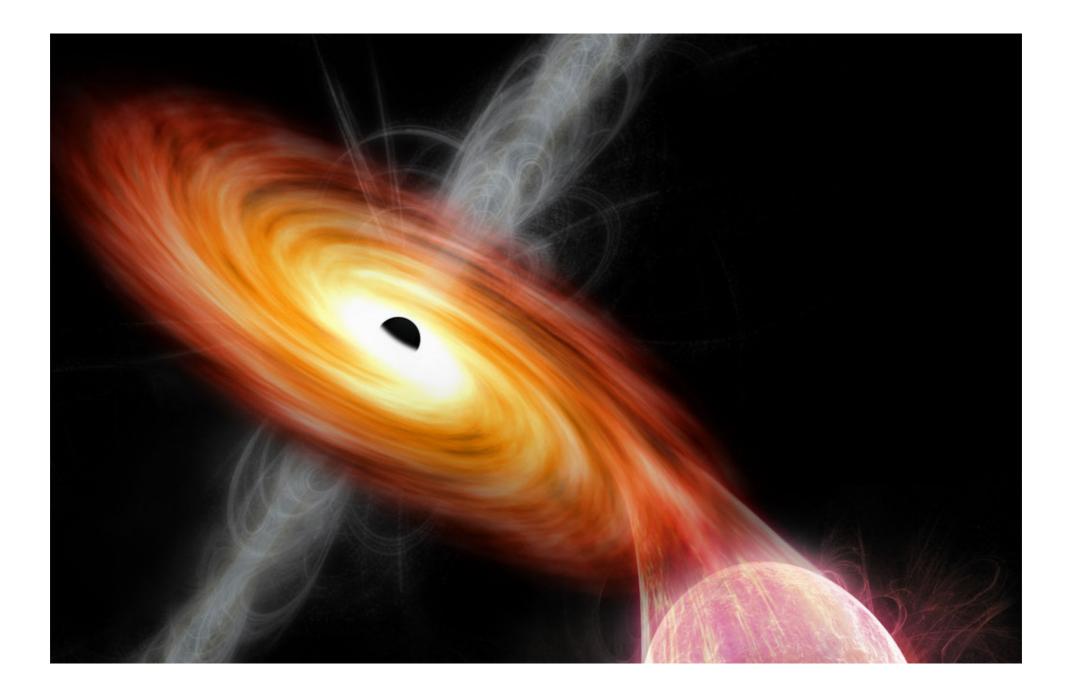
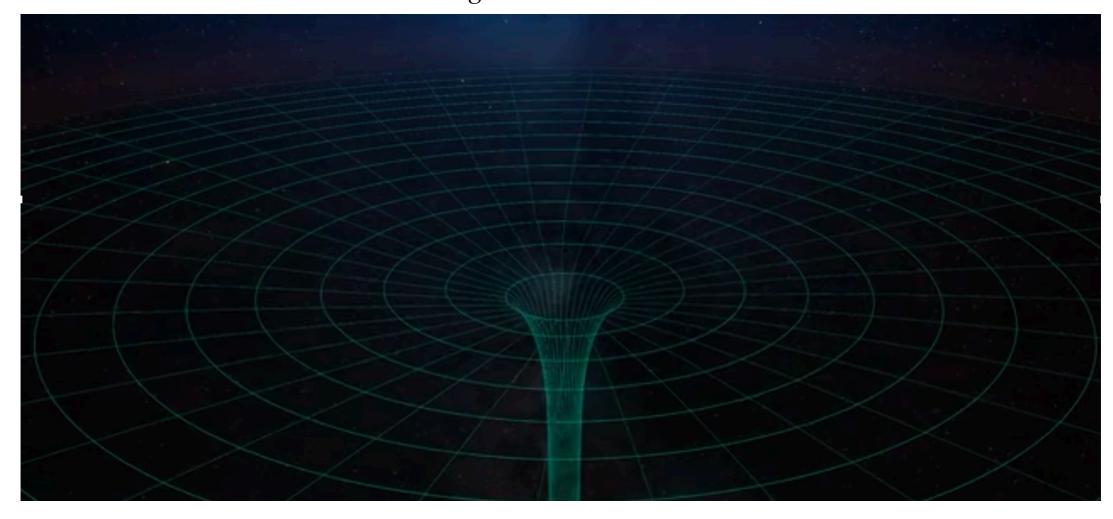
High Energy Astrophysics Dr. Adam Ingram



Lecture 4 Black Hole Accretion Discs

Black Holes

- All mass, M, in a singularity
- Event Horizon: $v_{esc} = c!$
- Newtonian approx: $v_{\rm esc}^2 = 2GM/r \implies r_h = 2GM/c^2$
- This is correct for a non-spinning black hole!
- Size scale: gravitational radius: $r_g = GM/c^2$

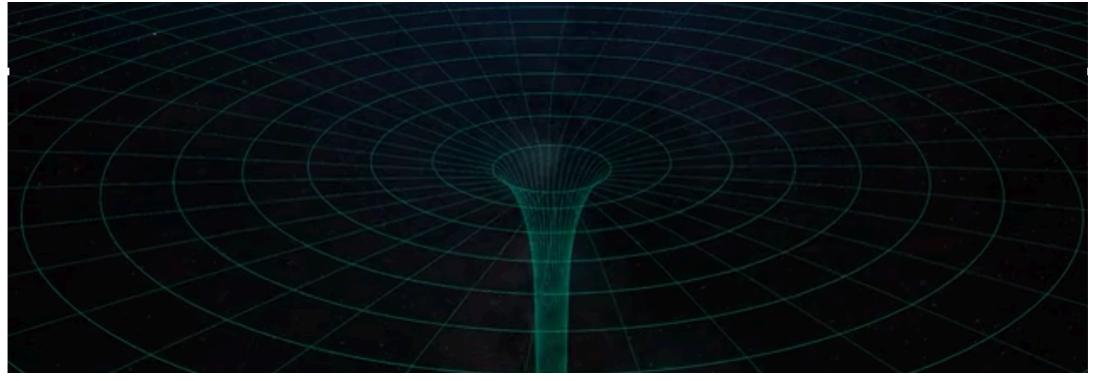


Non-spinning black hole: Spacetime described by Schwarzschild metric

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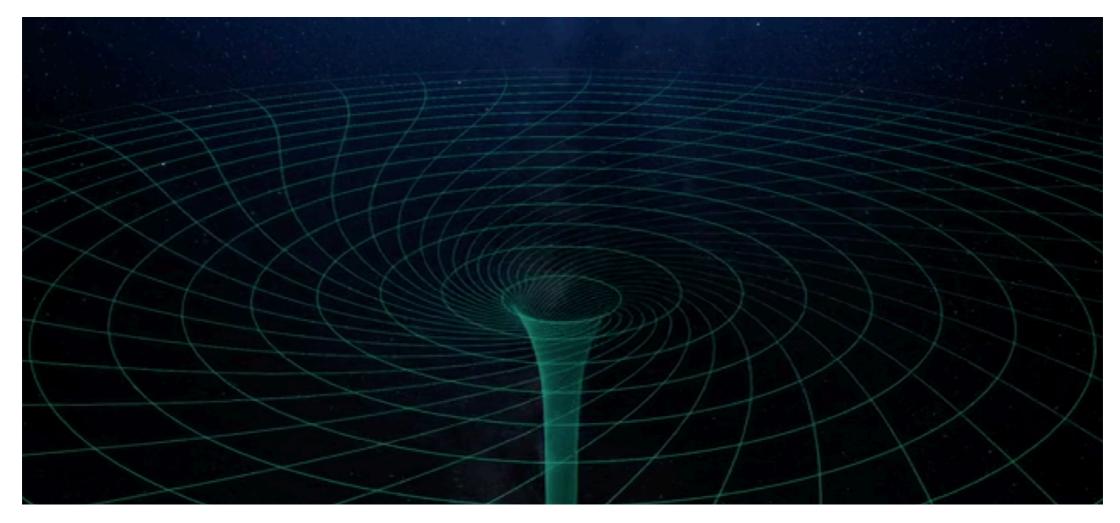
Symbol rg was gyroradius; is now gravitational radius!



Non-spinning black hole: Spacetime described by Schwarzschild metric

Black Holes

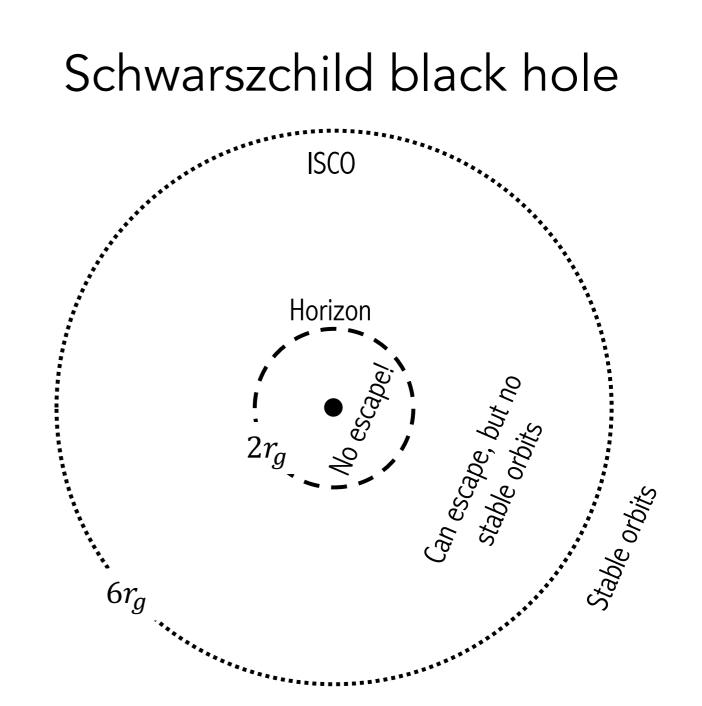
- Black hole can have angular momentum: $J_{\rm bh}$
- Dimensionless spin parameter: $a = J_{bh}/(Mcr_g)$
- Causality: $-1 \le a \le 1$...-ve for retrograde orbits.
- Horizon depends on spin: $r_h/r_g = 1 + \sqrt{1 a^2}$
- Rotating spacetime provides centrifugal barrier, so can get closer to BH



Non-spinning black hole: Spacetime described by Kerr metric

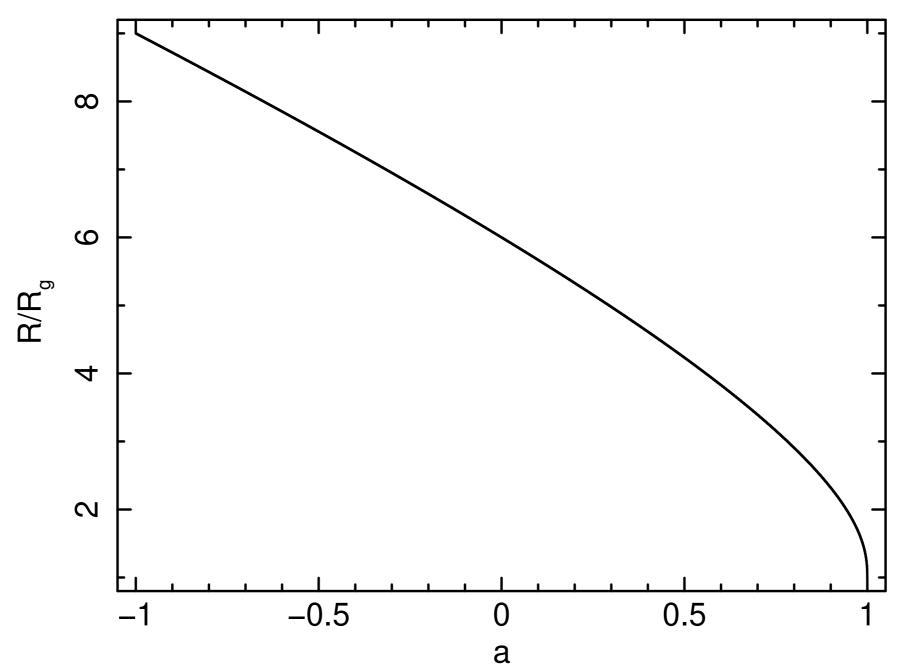
Innermost Stable Circular Orbit (ISCO)

- Stable circular orbits: gravitational attraction balanced by centrifugal barrier.
- Always have stable circular orbits in Newtonian gravity.
- In GR, close to a black hole gravity overcomes the centrifugal term: no stable orbits.



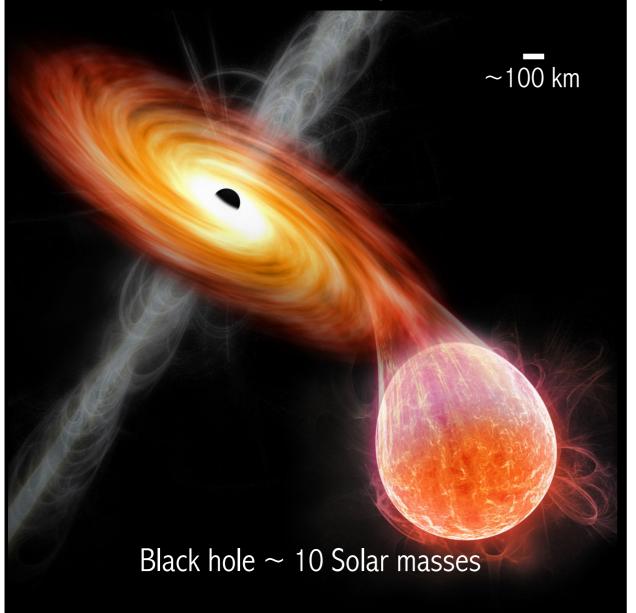
Innermost Stable Circular Orbit (ISCO)

ISCO depends on black hole spin: again, because spacetime is rotating (frame dragging effect), an orbiting test mass effectively has extra angular momentum, so black hole spin increases centrifugal barrier and enables stable orbits further in.



Accreting Black Holes

Black Hole X-ray Binaries



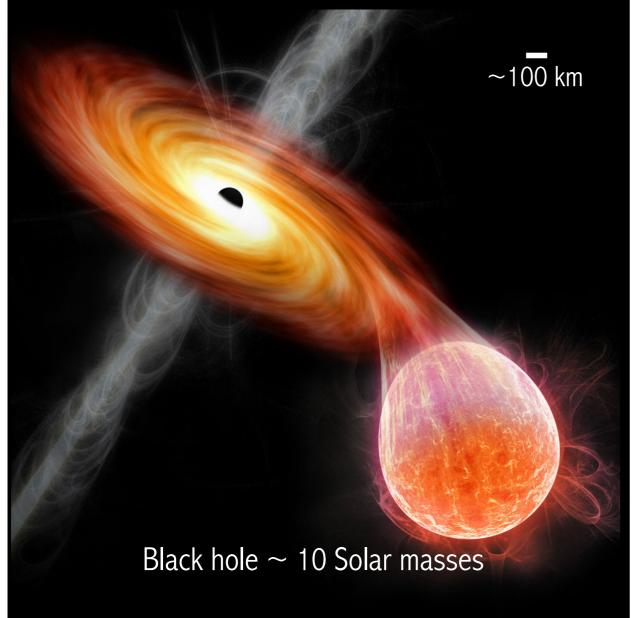
Active Galactic Nuclei



Black hole $\sim 10^{6}$ - 10^{10} Solar masses

Accreting Black Holes

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Active Galactic Nuclei



Black hole $\sim 10^{6}$ - 10^{10} Solar masses

Power supply: gravitational potential energy of accreting material. Therefore luminosity is the rest mass energy of accreted material multiplied by some efficiency factor: $L = \epsilon \dot{M} c^2$

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$$L_{\rm Edd} = \frac{4\pi GMcm_p}{\sigma_T} \qquad \qquad m_p = {\rm Proton\ mass} \\ \sigma_T = {\rm Thomson\ cross-section} \end{cases}$$

Assumptions: photon-electron scattering opacity, spherical accretion

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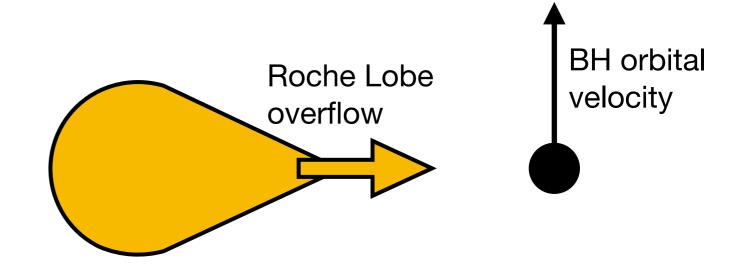
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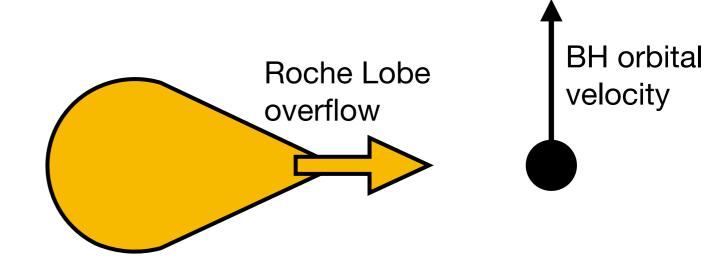
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- But what is the efficiency?

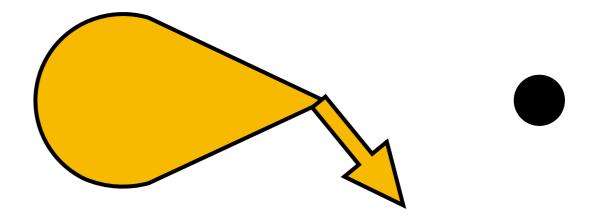
Companion star rest frame:



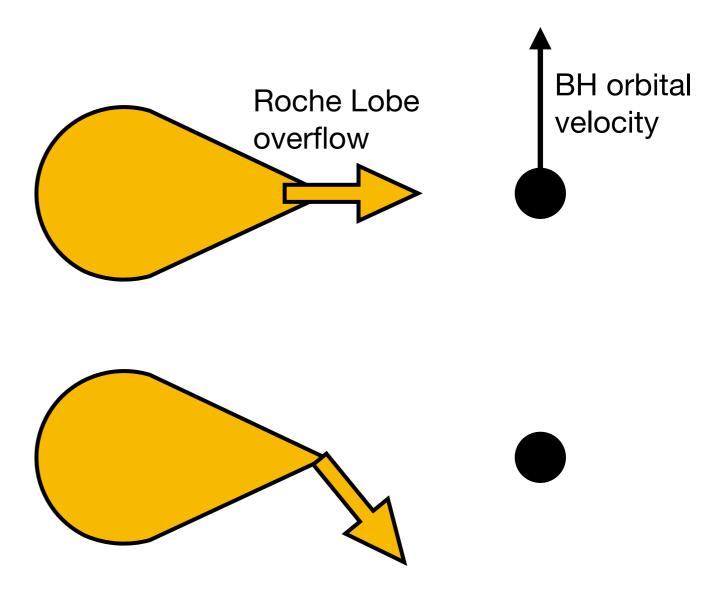
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Black hole rest frame:



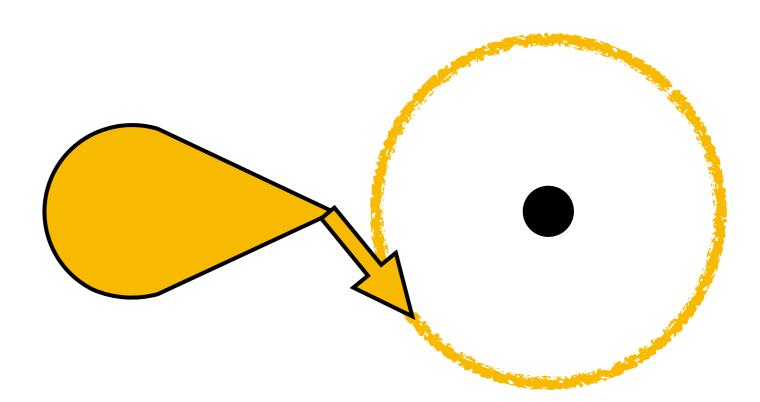
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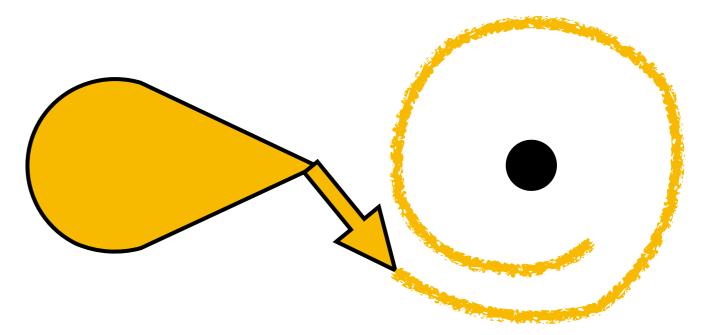
Black hole rest frame:

Therefore material has <u>angular momentum</u> about the black hole (similar arguments hold for AGN).

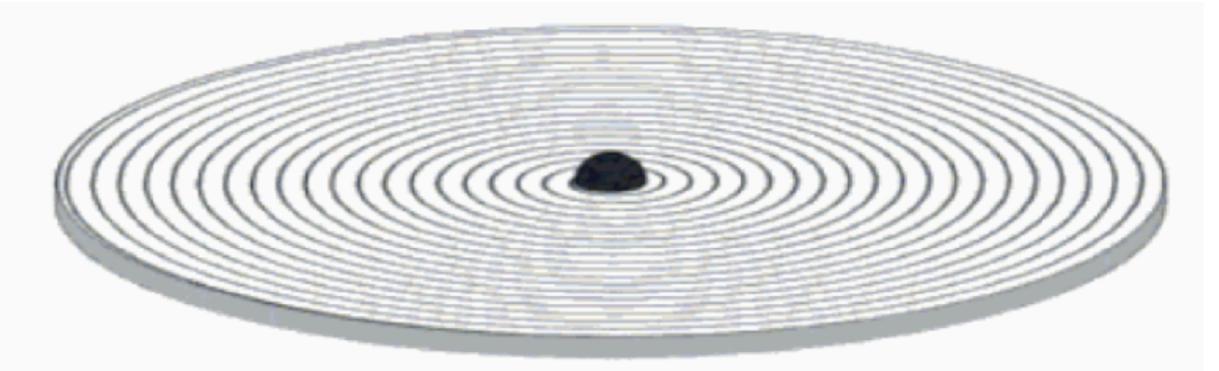
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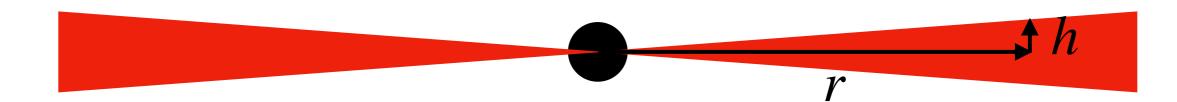
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• Disc scale height: $h/r \sim \text{constant}$



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• Therefore <u>half</u> of the GPE liberated can be radiated away (virial theorem). Luminosity of annulus is: $GM\dot{M}$

$$dL = \frac{GMM}{2r^2}dr$$

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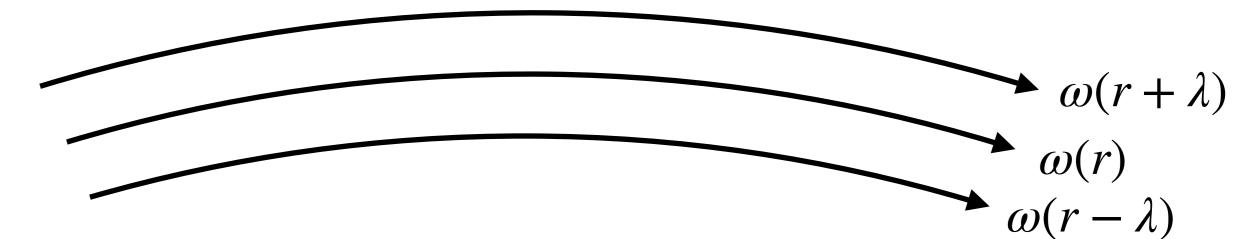
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Maximally spinning BH:

- $\epsilon = \frac{1}{12}$ $\epsilon = \frac{1}{2}$
- This is <u>enormous</u>! Nuclear fusion has efficiency $\epsilon \approx 0.007$

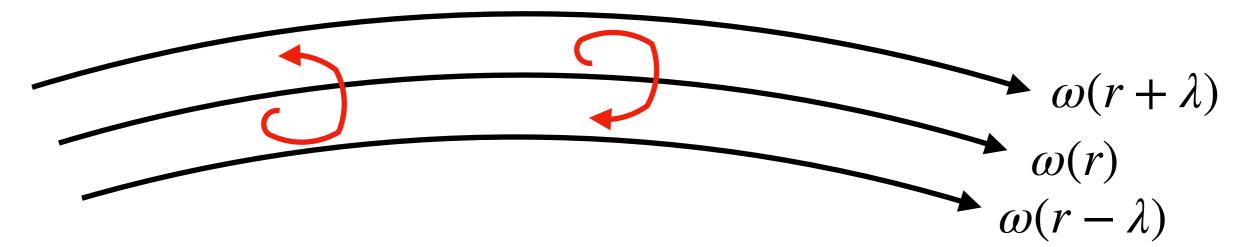
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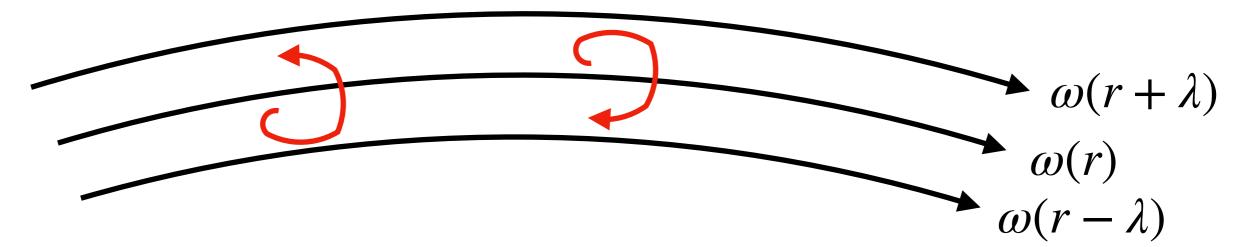
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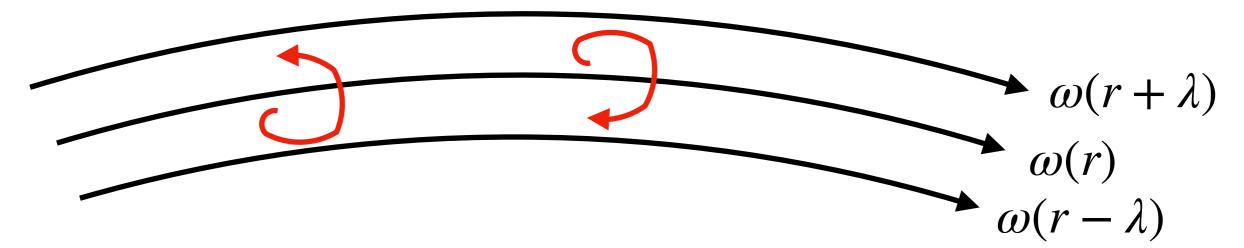
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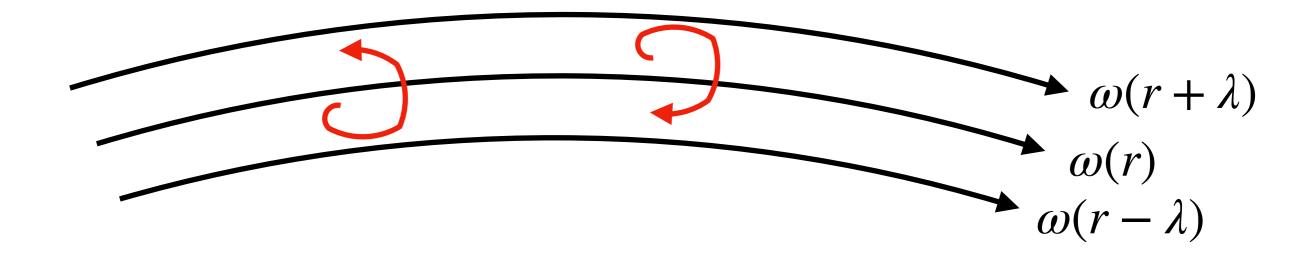
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- Therefore angular momentum transfer, since particles from the inner ring have larger angular velocity around BH than those from outer ring.
- In other words: viscosity (~friction) <u>slows down</u> the inner ring and <u>speeds up</u> the outer ring
- i.e. angular momentum transported outwards, material can spiral towards the BH.

So what is causing the turbulent motions (~viscosity)?

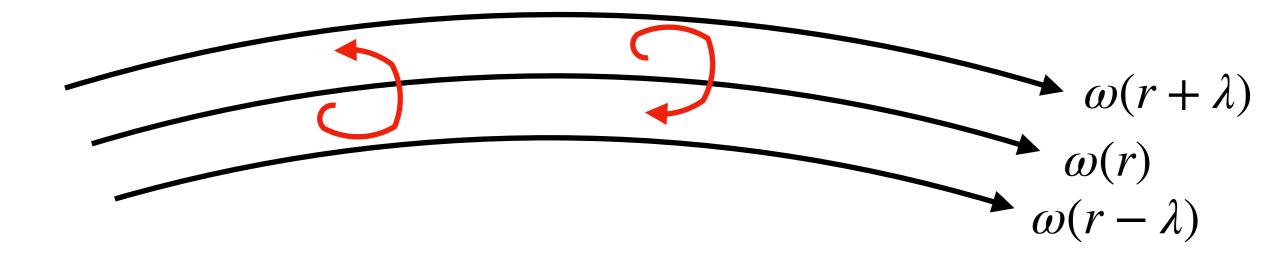


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• For viscosity due to random thermal motions of particles:

 $\bar{v} \sim c_s$ = sound speed $\lambda \sim$ mean free path of particles

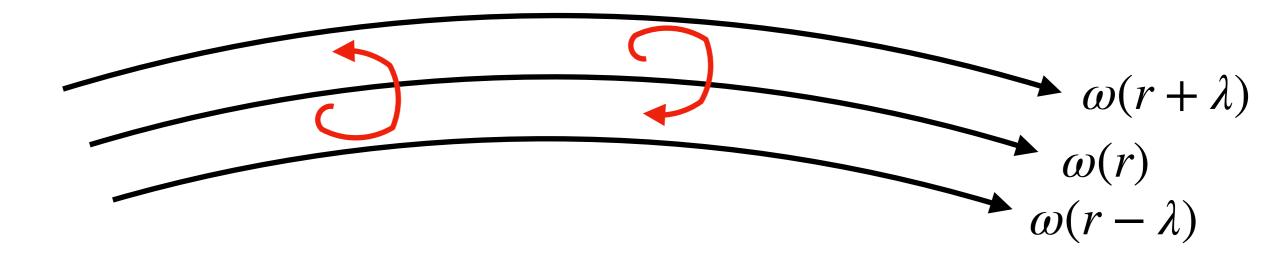
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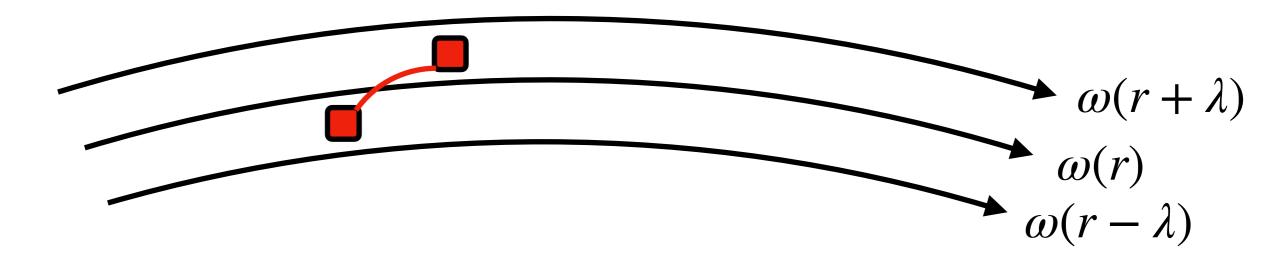
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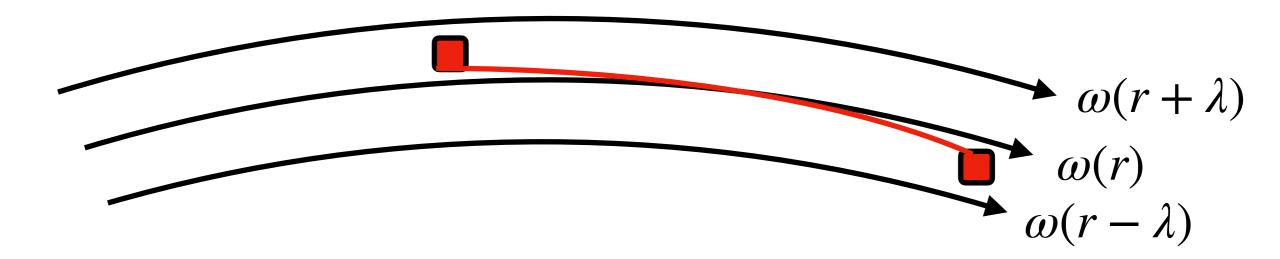
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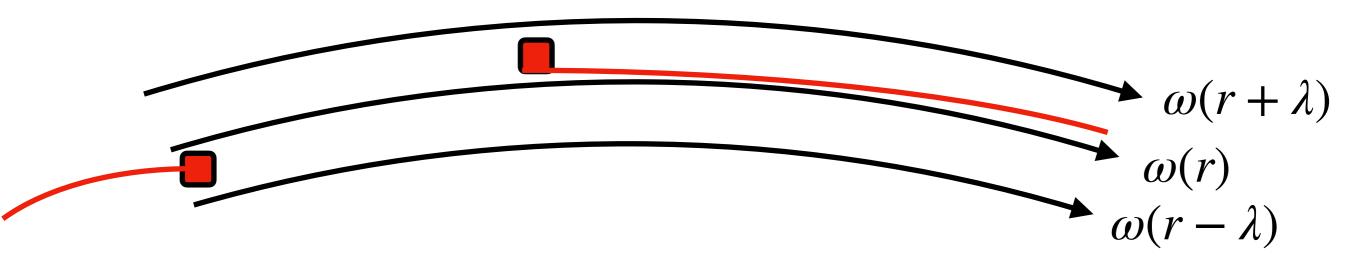
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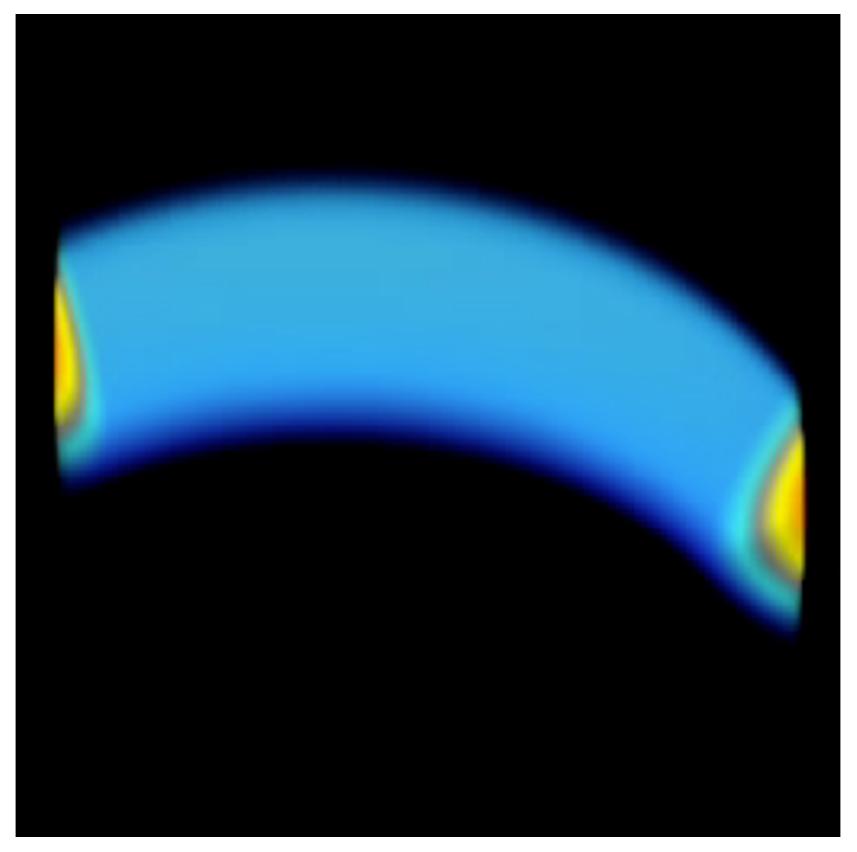


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- Effective viscosity must be provided by *hydrodynamic turbulence*.
- The magneto-rotational instability provides the required viscosity: B-field lines connect parcels of gas, differential rotation stretches distance between these parcels, causes field lines to become tangled and generate turbulence. Like winding an elastic band round and round the disc.





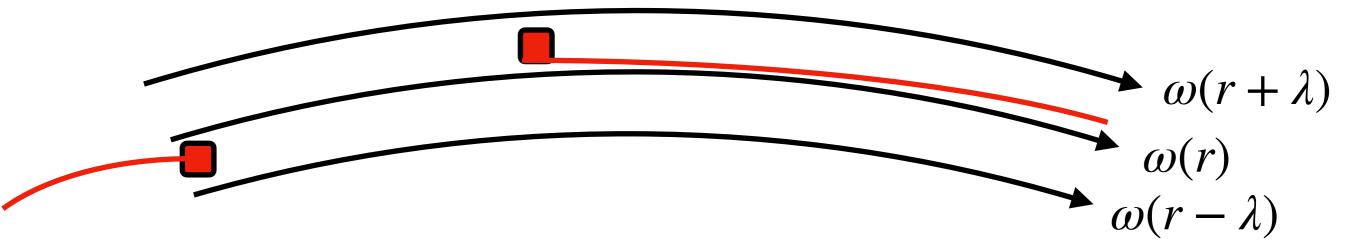
Krolik, de Villiers & Hawley

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• Shakura & Sunyaev (1973) famously assumed for the *kinematic viscosity*:

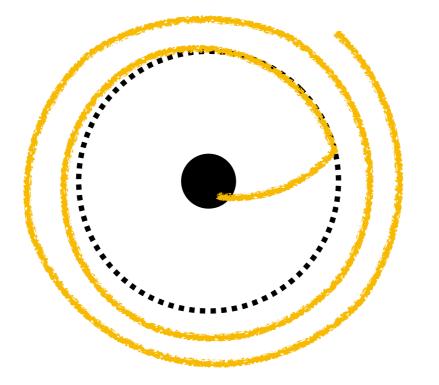
$$\nu = \alpha c_s h$$

- α is a dimensionless constant: speed of eddies $\lesssim c_s$; length scale of eddies $\lesssim h$, therefore $\alpha \lesssim 1$.
- The "alpha-disc" model has been remarkably successful.



What happens at r_{isco}?

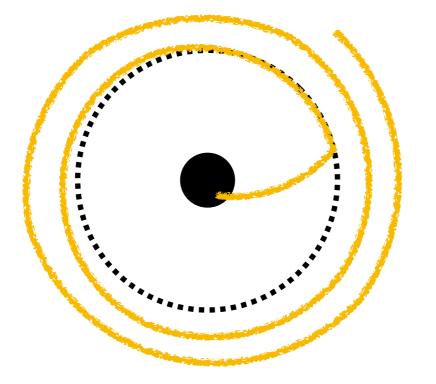
• Inside of ISCO, no stable orbits so material is in free fall.



- Often assumed there is therefore no "stress" at ISCO (i.e. viscosity parameter drops to zero).
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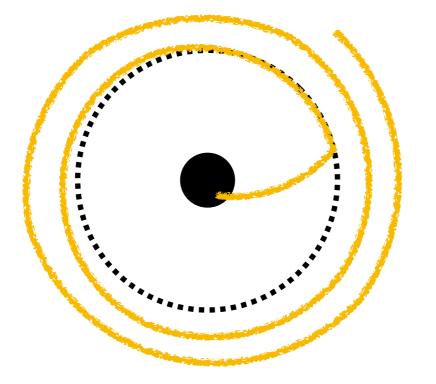


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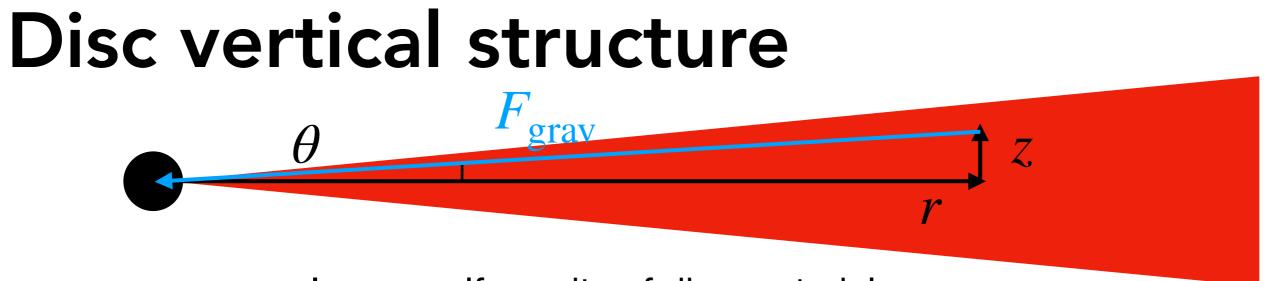
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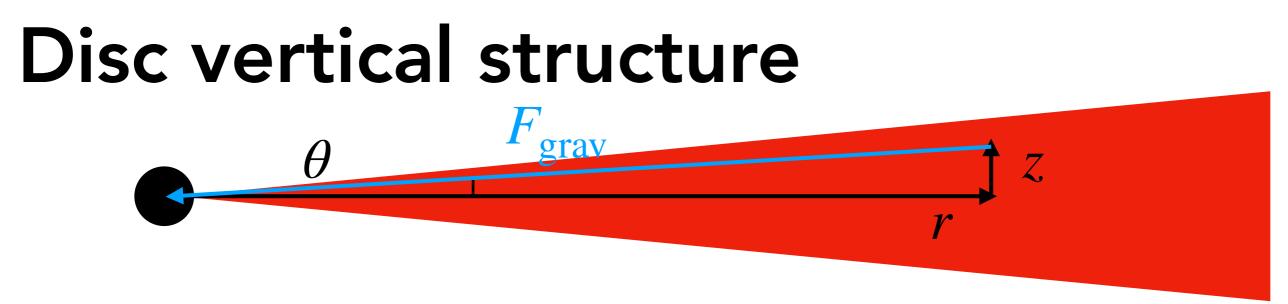
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... It increases the mass of the black hole!

• Therefore lower spin supermassive black holes can presumably grow faster than higher spin ones.



Ignore self-gravity of disc material

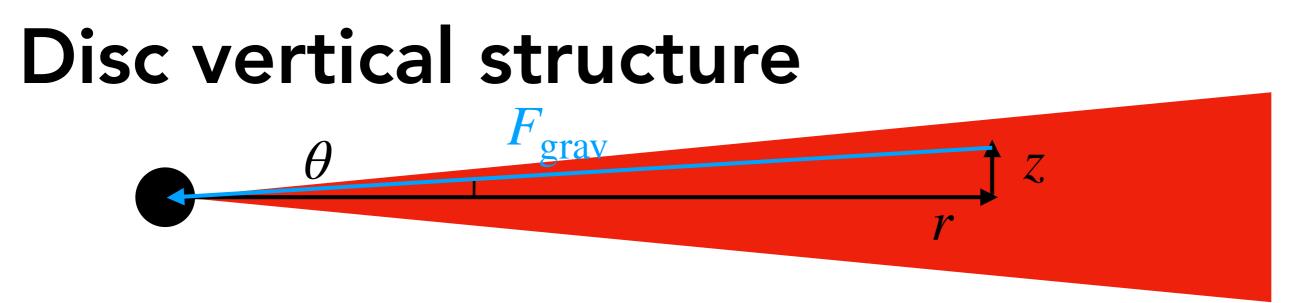


dP

dz

 $-\rho g_z$

 g_z is vertical component of gravitational acceleration, ρ is mass density



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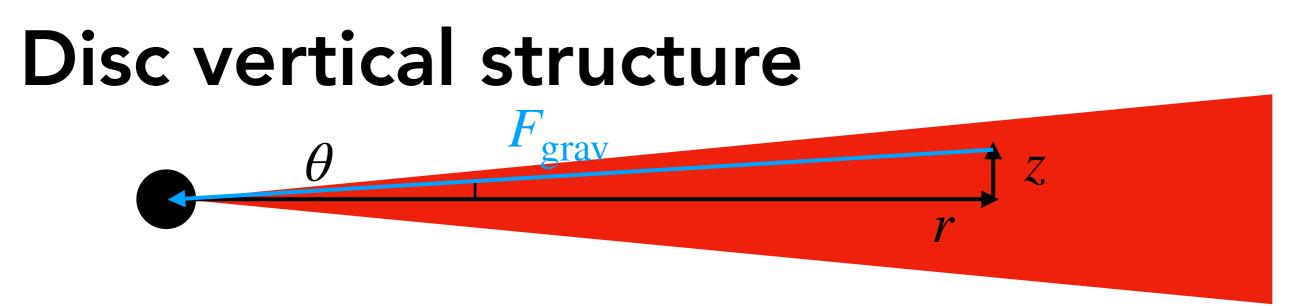
• Vertical component of gravitational acceleration is:

$$g_z = \frac{GM}{r^2 + z^2} \sin \theta$$

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dP

dz



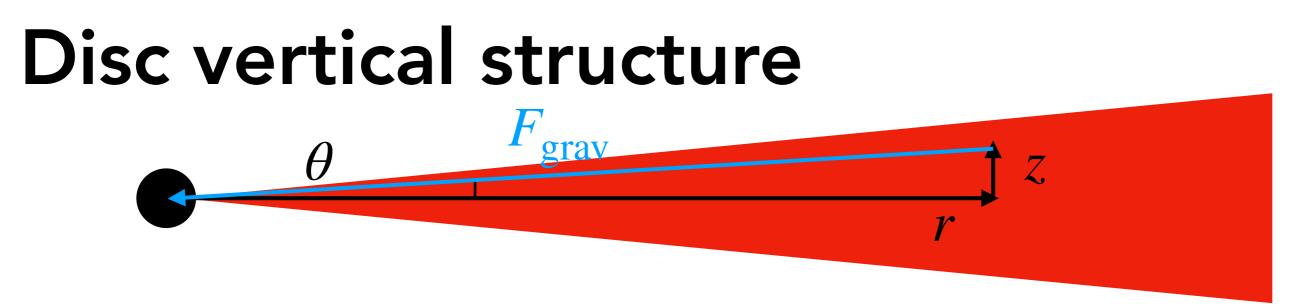
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(small angle approximation)



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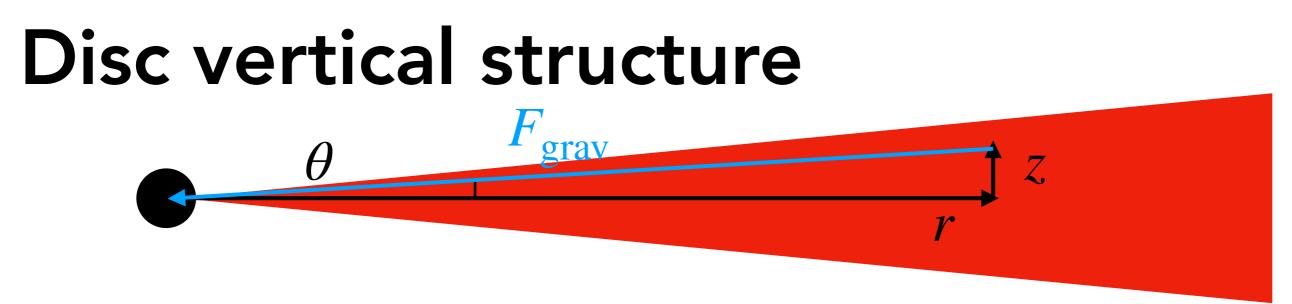
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• Pressure and density related via sound speed c_s:

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• Therefore vertical density gradient is:

$$\frac{d\rho}{dz} = -\rho \frac{GMz}{r^3 c_s^2}$$

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• Integrate to get: $\rho(r) = \rho_0 \exp\left[-\frac{GM}{c_s^2 r^3} \frac{z^2}{2}\right]$

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• Integrate to get:

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• Therefore density is Gaussian with peak at z=0 and width h:

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$$\frac{h}{r} \ll 1 \implies \underline{disc rotation is highly supersonic.}$$

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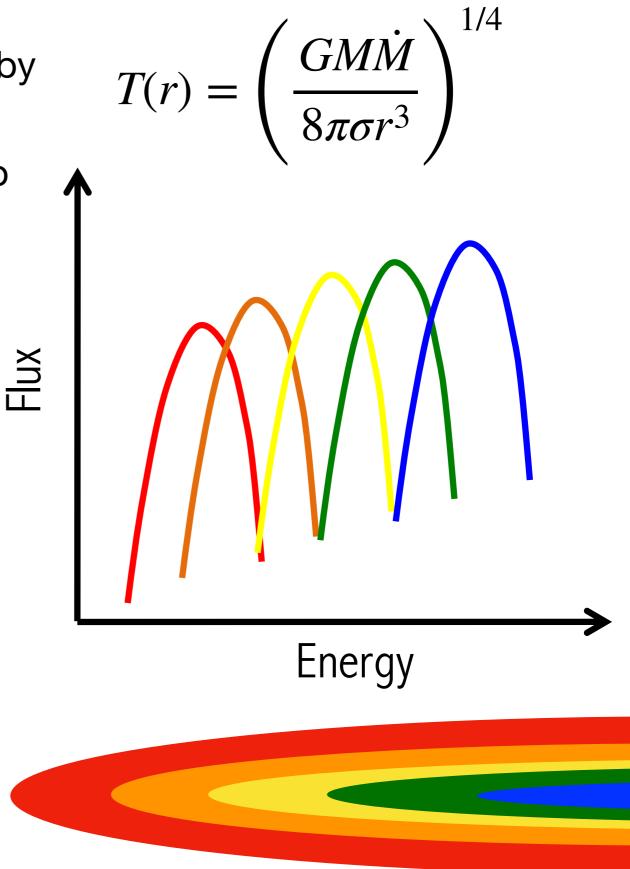
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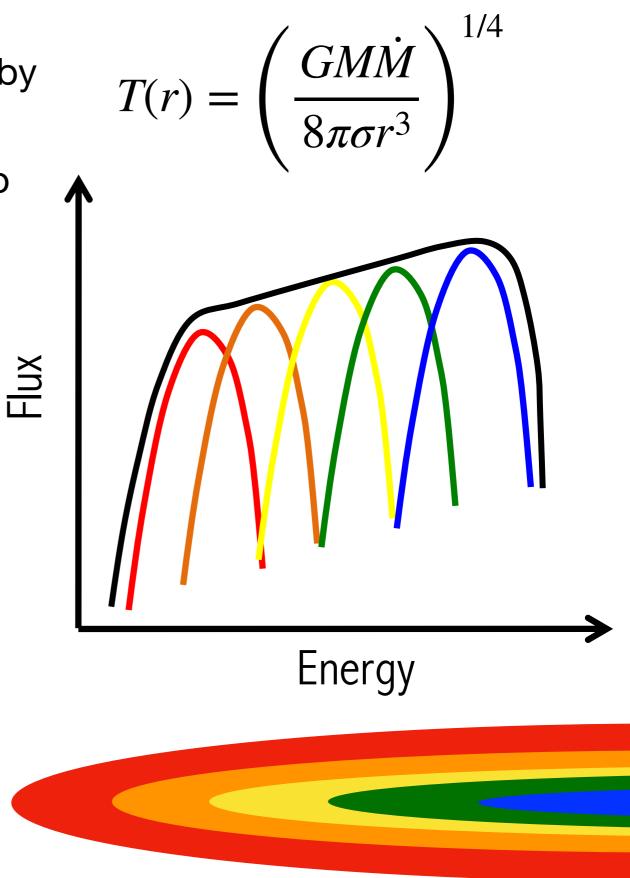
• From Stefan-Boltzmann law, the temperature is:

$$T(r) = \left(\frac{GM\dot{M}}{8\pi\sigma r^3}\right)^{1/4}$$

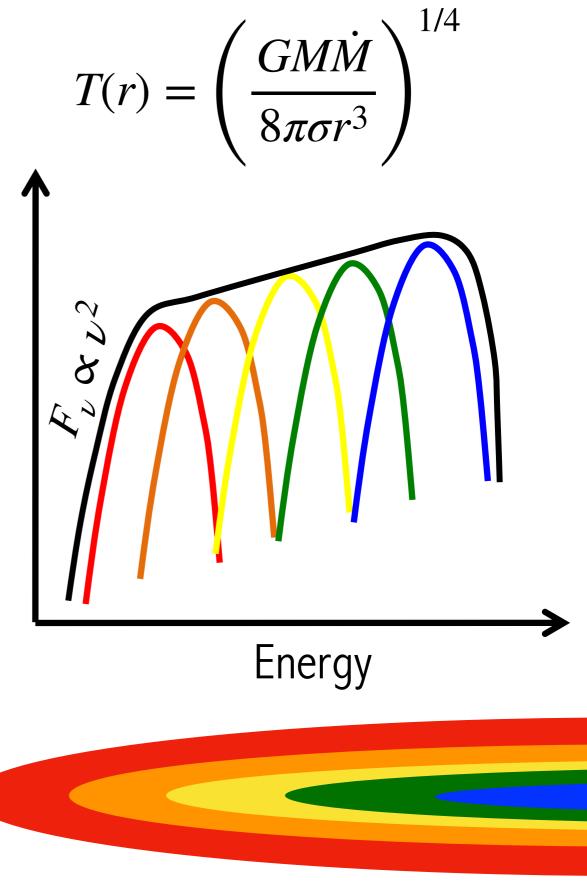
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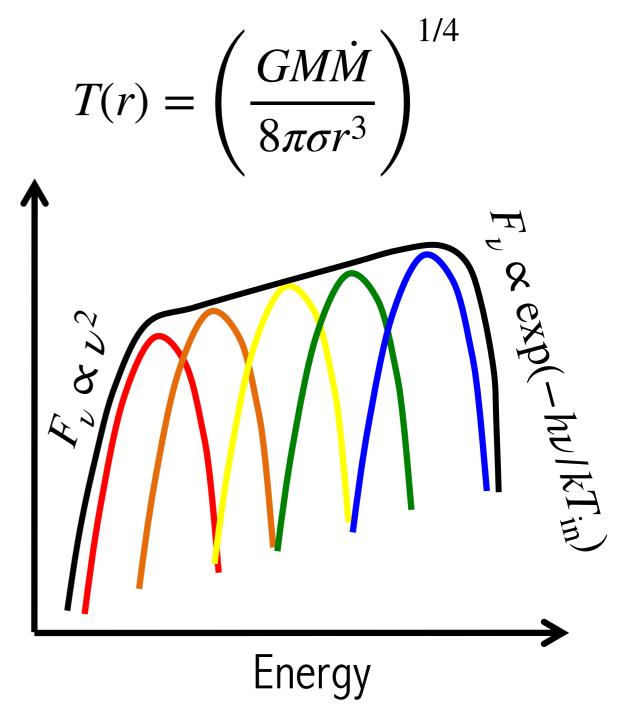
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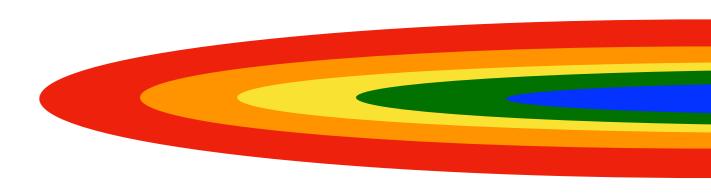


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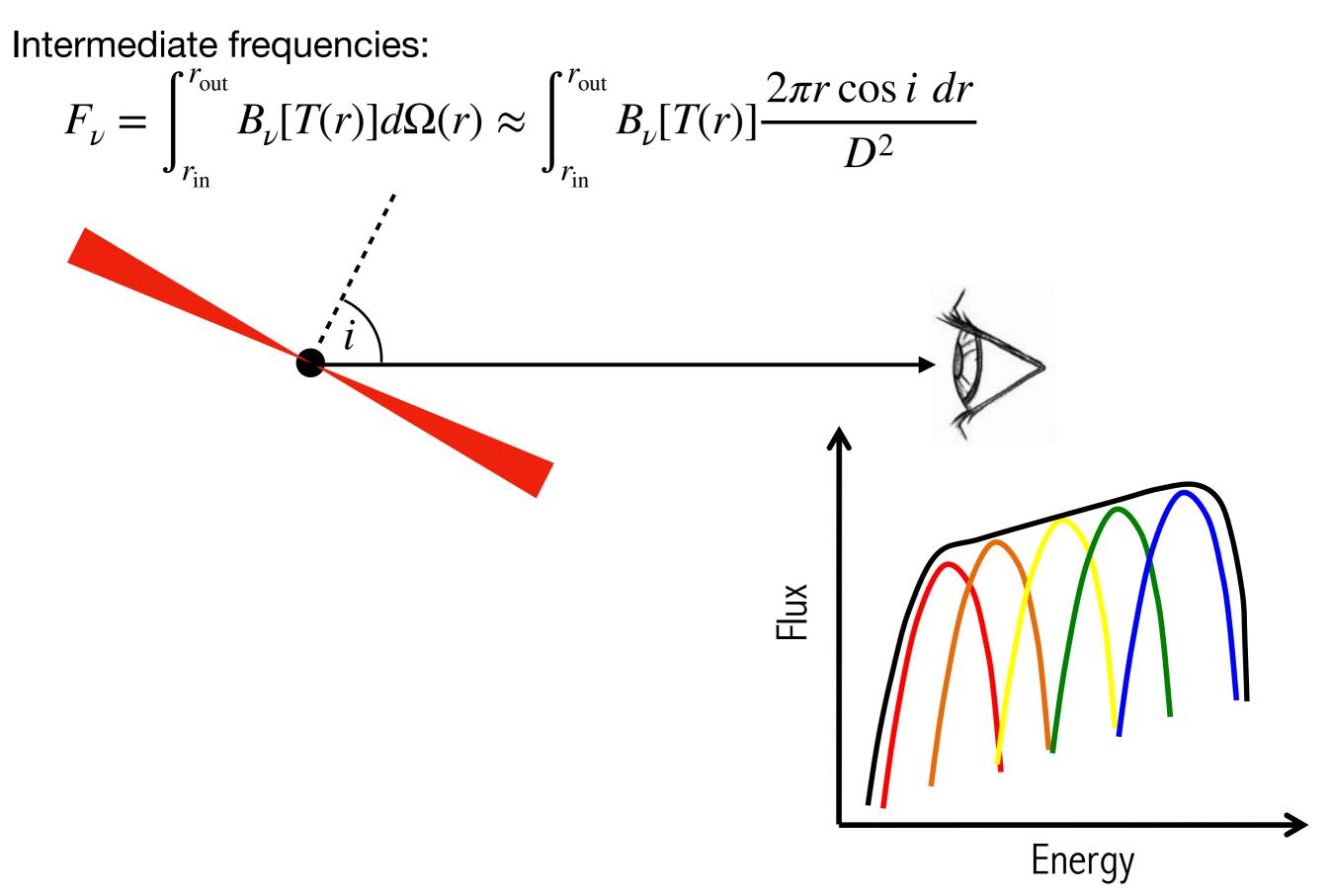


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Disc Spectrum

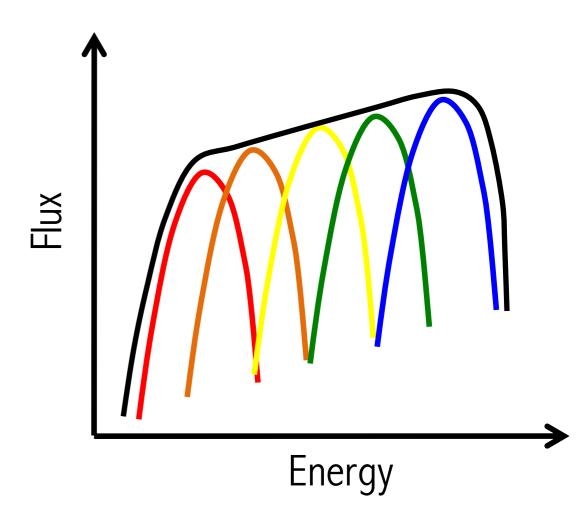


Disc Spectrum

Intermediate frequencies:

$$F_{\nu} \propto \int_{r_{\rm in}}^{r_{\rm out}} B_{\nu}[T(r)] r dr \propto \int_{r_{\rm in}}^{r_{\rm out}} \frac{1}{\exp(R)}$$

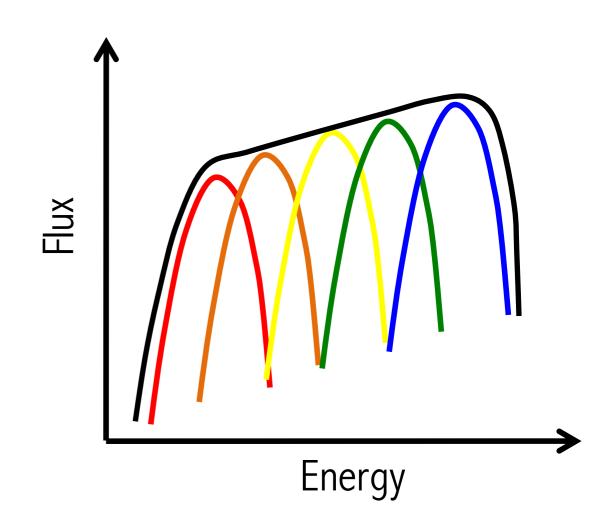
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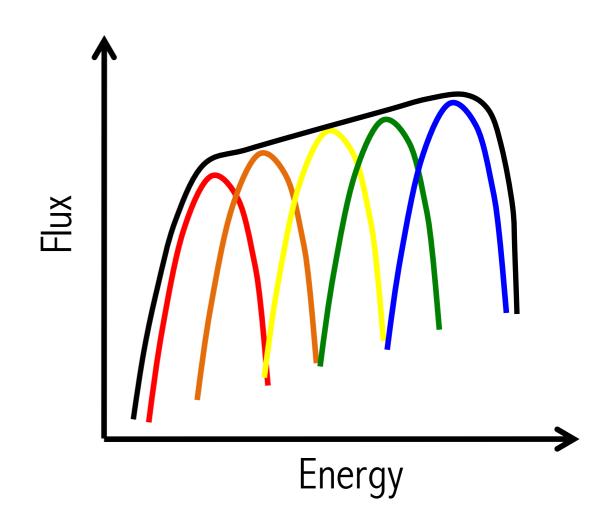
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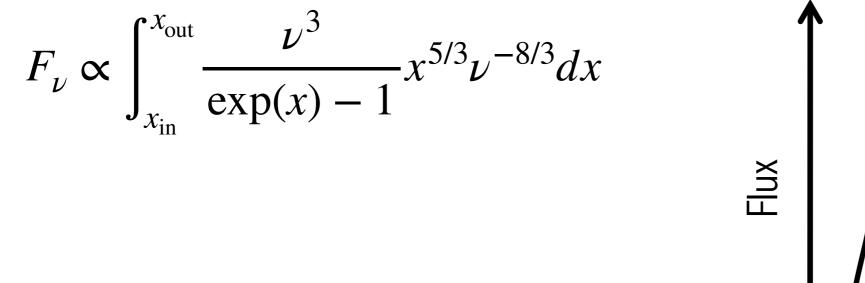


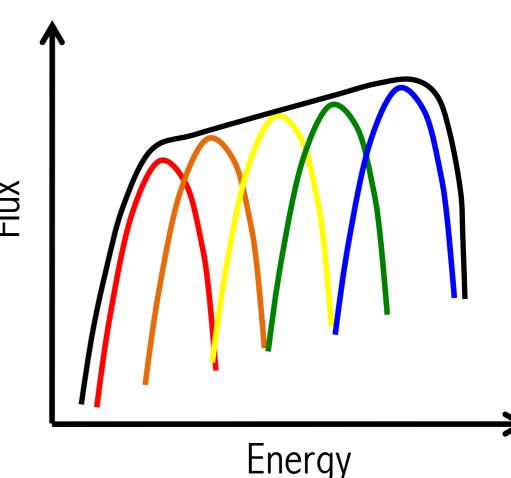
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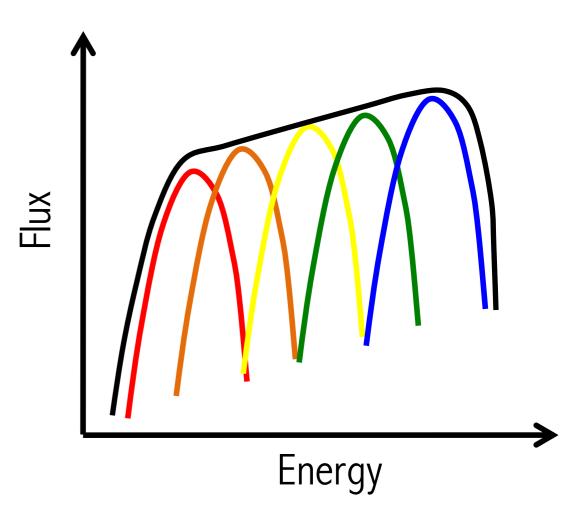
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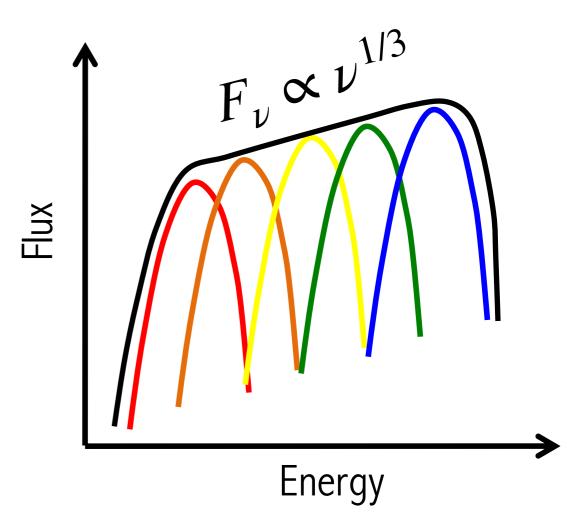
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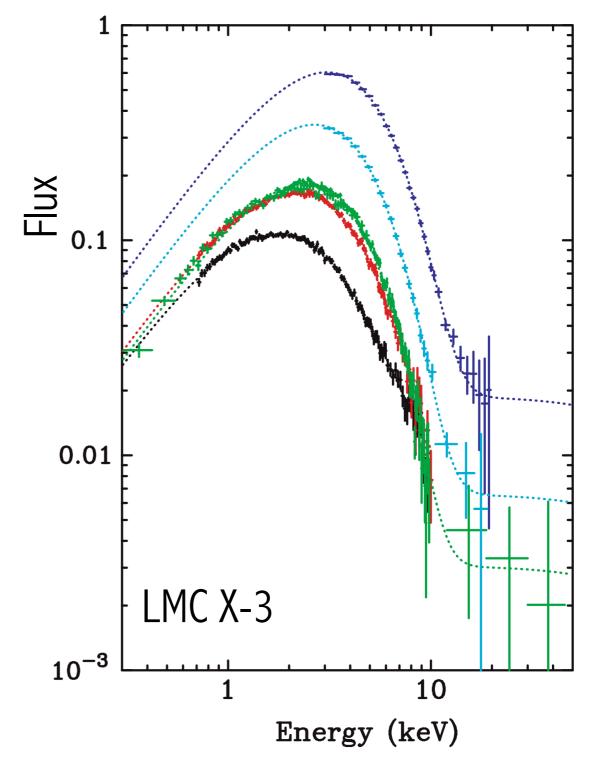
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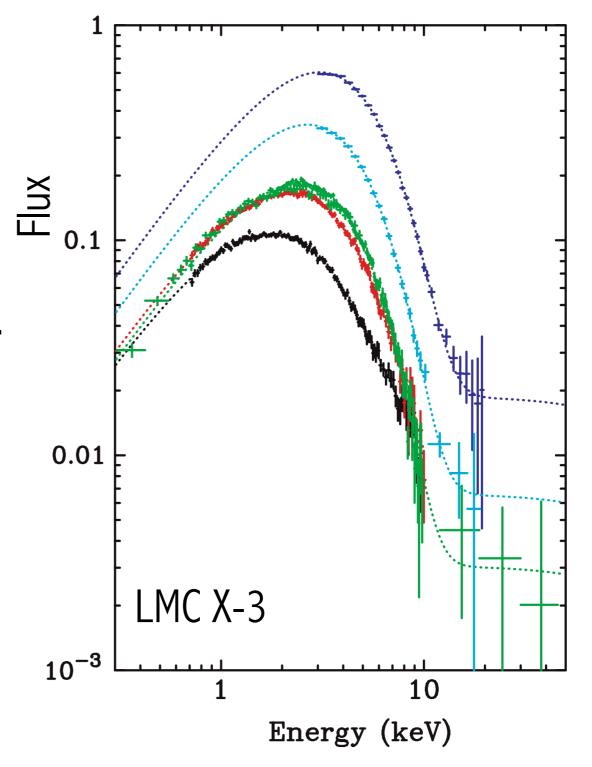


Do we see this disc spectrum from black hole X-ray binaries?

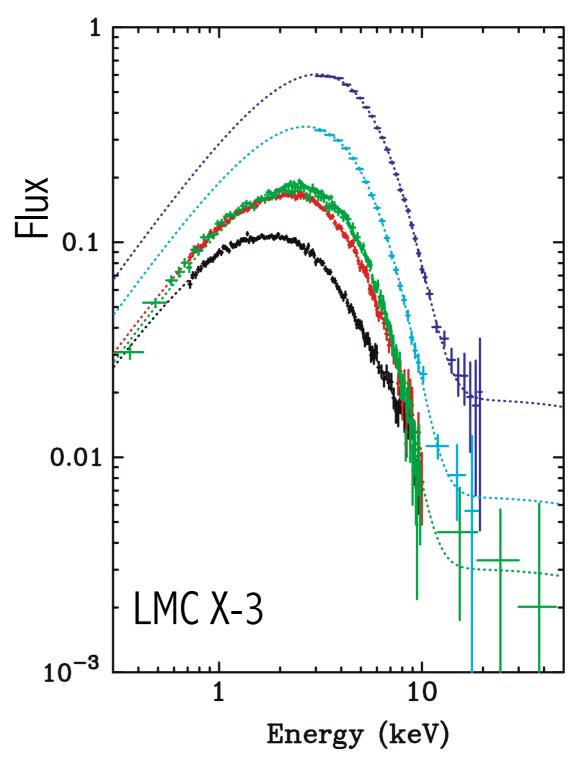
• Yes!



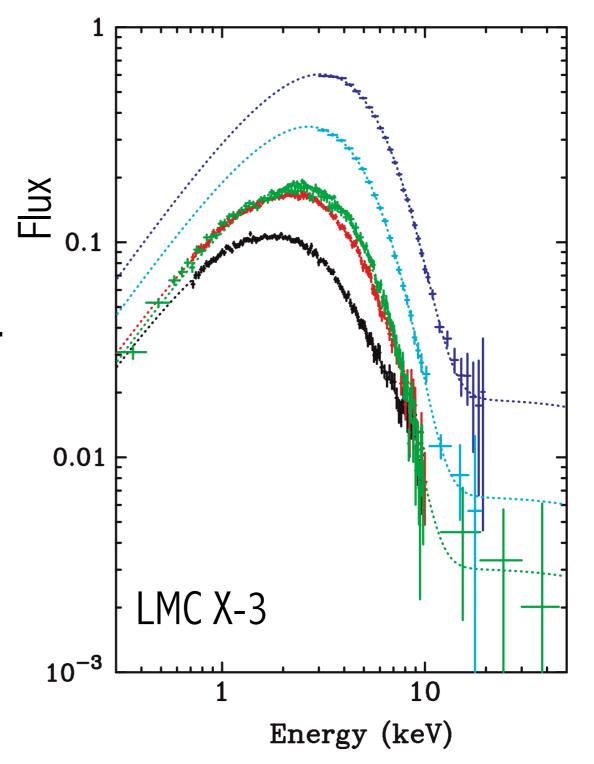
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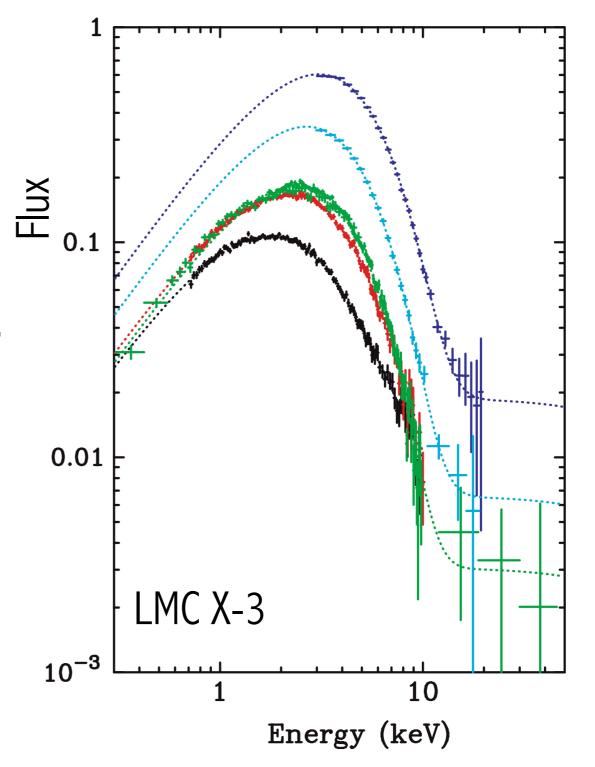
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- Find a range a~0.2-0.99 for different stellar-mass black holes

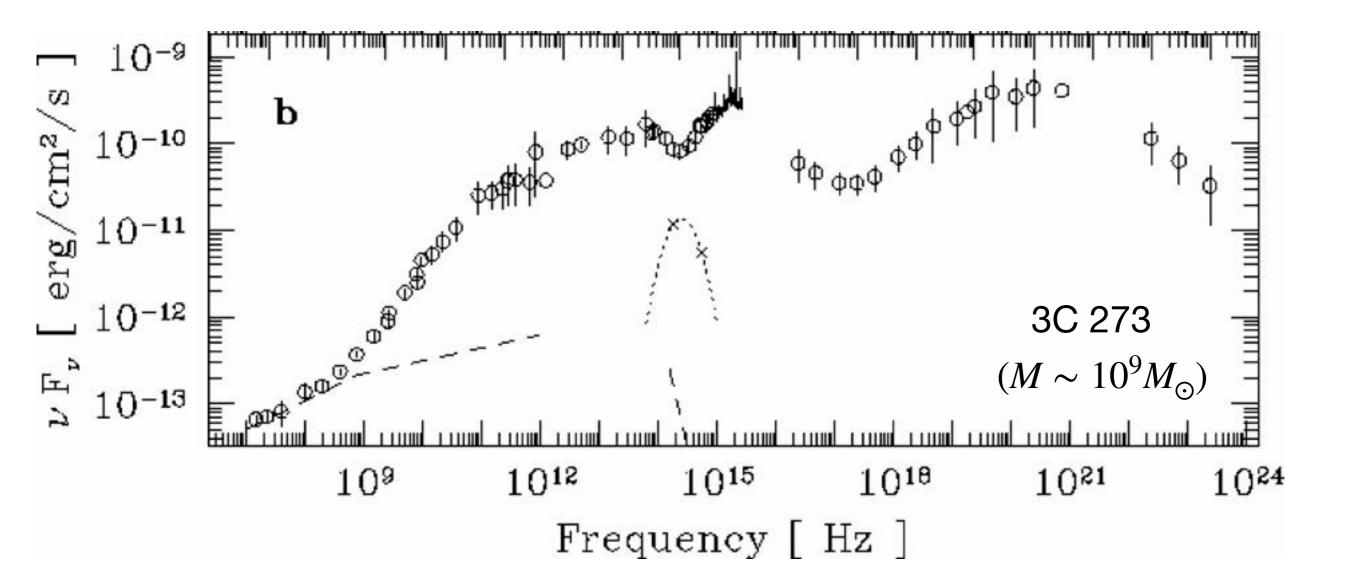
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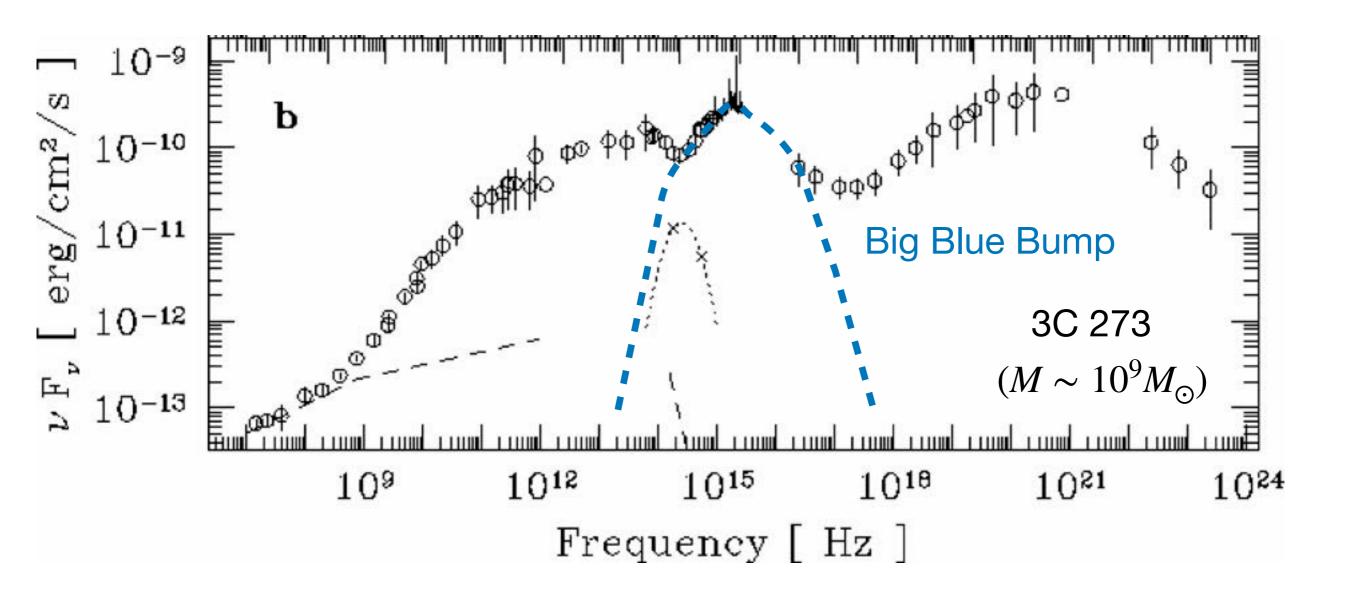
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- See "big blue bump" at expected frequency range.
- But see lots of other stuff, some of which we will study over the following two lectures.

