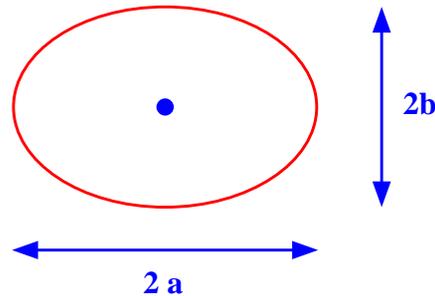


7. BINARY STARS (ZG: 12; CO: 7, 17)

- most stars are members of binary systems or multiple systems (triples, quadruples, quintuplets, ...)
- *orbital period* distribution: $P_{\text{orb}} = 11 \text{ min to } \sim 10^6 \text{ yr}$
- the majority of binaries are wide and do not interact strongly
- *close binaries* (with $P_{\text{orb}} \lesssim 10 \text{ yr}$) can transfer mass \rightarrow *changes structure and subsequent evolution*
- *approximate period distribution*: $f(\log P) \simeq \text{const.}$
(rule of thumb: 10% of systems in each decade of $\log P$ from 10^{-3} to 10^7 yr)

generally large scatter in distribution of eccentricities

- $e^2 \equiv 1 - b^2/a^2$,
a = semi-major,
b = semi-minor axis



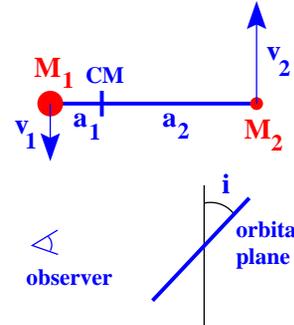
- systems with eccentricities $\lesssim 10$ tend to be circular \rightarrow evidence for *tidal circularization*

7.1 Classification

- *visual binaries*: see the periodic wobbling of two stars in the sky (e.g. Sirius A and B); if the motion of only one star is seen: *astrometric binary*
- *spectroscopic binaries*: see the periodic *Doppler shifts* of spectral lines
 - ▷ *single-lined*: only the Doppler shifts of one star detected
 - ▷ *double-lined*: lines of both stars are detected
- *photometric binaries*: periodic variation of fluxes, colours, etc. are observed (caveat: such variations can also be caused by single variable stars: Cepheids, RR Lyrae variables)
- *eclipsing binaries*: one or both stars are eclipsed by the other one \rightarrow inclination of orbital plane $i \simeq 90^\circ$ (most useful for determining basic stellar parameters)

7.2 THE BINARY MASS FUNCTION

- consider a *spectroscopic binary*
- measure the *radial velocity curve* along the line of sight from $\frac{v_r}{c} \simeq \frac{\Delta\lambda}{\lambda}$ (Doppler shift)



$$\triangleright M_1 a_1 = M_2 a_2$$

$$\triangleright P = \frac{2\pi}{\omega} = 2\pi \frac{a_1 \sin i}{v_1 \sin i} = 2\pi \frac{a_2 \sin i}{v_2 \sin i}$$

$$\triangleright \text{gravitational force} = \text{centripetal force}$$

$$\rightarrow \frac{GM_1 M_2}{(a_1 + a_2)^2} = \frac{(v_1 \sin i)^2}{a_1 \sin^2 i} M_1, \quad \frac{GM_1 M_2}{(a_1 + a_2)^2} = \frac{(v_2 \sin i)^2}{a_2 \sin^2 i} M_2$$

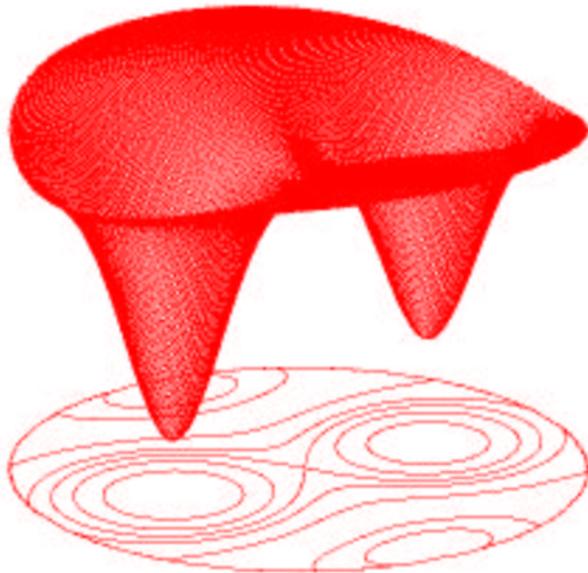
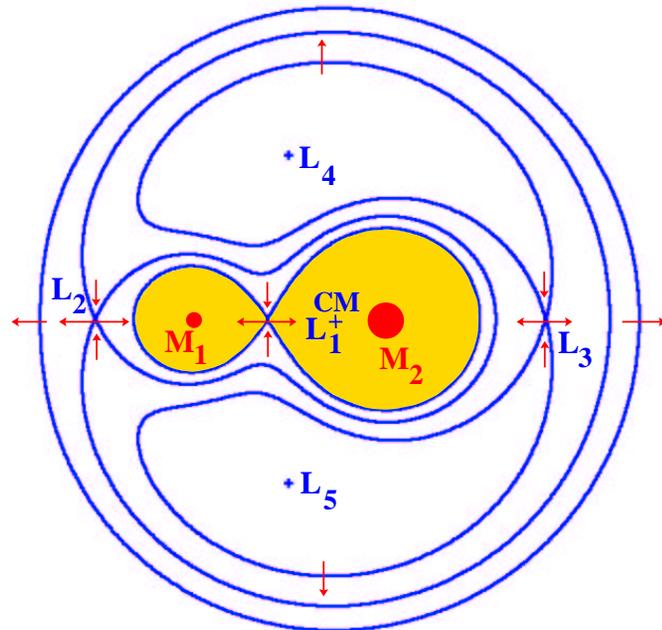
substituting $(a_1 + a_2)^2 = a_1^2 (M_1 + M_2)^2 / M_2^2$, etc.

$$\rightarrow f_1(M_2) = \frac{M_2^3 \sin^3 i}{(M_1 + M_2)^2} = \frac{P (v_1 \sin i)^3}{2\pi G}$$

$$f_2(M_1) = \frac{M_1^3 \sin^3 i}{(M_1 + M_2)^2} = \frac{P (v_2 \sin i)^3}{2\pi G}$$

- f_1, f_2 *mass functions*: relate observables $v_1 \sin i, v_2 \sin i, P$ to quantities of interest $M_1, M_2, \sin i$
- measurement of f_1 and f_2 (for double-lined spectroscopic binaries only) determines $M_1 \sin^3 i, M_2 \sin^3 i$
 - \triangleright if i is known (e.g. for visual binaries or eclipsing binaries) $\rightarrow M_1, M_2$
 - \triangleright for $M_1 \ll M_2 \rightarrow f_1(M_2) \simeq M_2 \sin^3 i$ (measuring $v_1 \sin i$ for star 1 constrains M_2)
- for *eclipsing binaries* one can also determine the *radii* of both stars (main source of accurate masses and radii of stars [and luminosities if distances are known])

The Roche Potential



7.3 THE ROCHE POTENTIAL

- *restricted three-body problem*: determine the motion of a test particle in the field of two masses M_1 and M_2 in a circular orbit about each other
- equation of motion of the particle in a frame rotating with the binary $\Omega = 2\pi/P$:

$$\frac{d^2\vec{r}}{dt^2} = -\vec{\nabla} U_{\text{eff}} - \underbrace{\frac{2\vec{\Omega} \times \vec{v}}{\text{Coriolis force}}}$$

where the *effective potential* U_{eff} is given by

$$U_{\text{eff}} = -\frac{GM_1}{|\vec{r} - \vec{r}_1|} - \frac{GM_2}{|\vec{r} - \vec{r}_2|} - \underbrace{\frac{1}{2}\Omega^2(x^2 + y^2)}_{\text{centrifugal term}}$$

- *Lagrangian points*: five stationary points of the Roche potential U_{eff} (i.e. where effective gravity $\vec{\nabla}U_{\text{eff}} = 0$)
 - ▷ 3 saddle points: L_1, L_2, L_3
- *Roche lobe*: equipotential surface passing through the *inner Lagrangian point* L_1 ('connects' the gravitational fields of the two stars)
- approximate formula for the *effective Roche-lobe radius* (of star 2):

$$R_L = \frac{0.49}{0.6 + q^{2/3} \ln(1 + q^{-1/3})} A,$$

where $q = M_1/M_2$ is the mass ratio, A orbital separation.

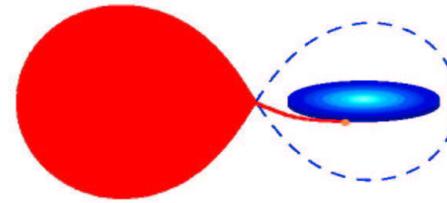
Classification of close binaries

- **Detached binaries:**
 - ▷ both stars *underfill their Roche lobes*, i.e. the photospheres of both stars lie beneath their respective Roche lobes
 - ▷ *gravitational interactions* only (e.g. tidal interaction, see Earth-Moon system)
- **Semidetached binaries:**
 - ▷ one star *fills its Roche lobe*
 - ▷ the Roche-lobe filling component *transfers matter* to the detached component
 - ▷ *mass-transferring binaries*
- **Contact binaries:**
 - ▷ *both stars fill or overflow their Roche lobes*
 - ▷ formation of a common photosphere surrounding both components
 - ▷ *W Ursae Majoris stars*

7.4 BINARY MASS TRANSFER

- 30 - 50 % of all stars experience mass transfer by *Roche-lobe overflow* during their lifetimes (generally in late evolutionary phases)

a) (quasi-)conservative mass transfer



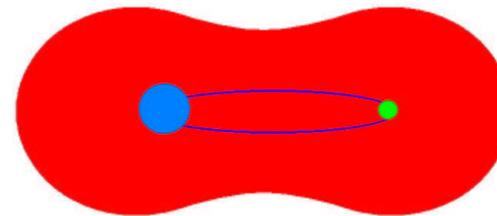
▷ *mass loss + mass accretion*

▷ the mass loser tends to lose most of its envelope → formation of *helium stars*

▷ the accretor tends to be *rejuvenated* (i.e. behaves like a more massive star with the evolutionary clock reset)

▷ *orbit generally widens*

b) dynamical mass transfer → common-envelope and spiral-in phase (mass loser is usually a red giant)



▷ accreting component also fills its Roche lobe

▷ mass donor (primary) *engulfs secondary*

▷ *spiral-in* of the core of the primary and the secondary immersed in a *common envelope*

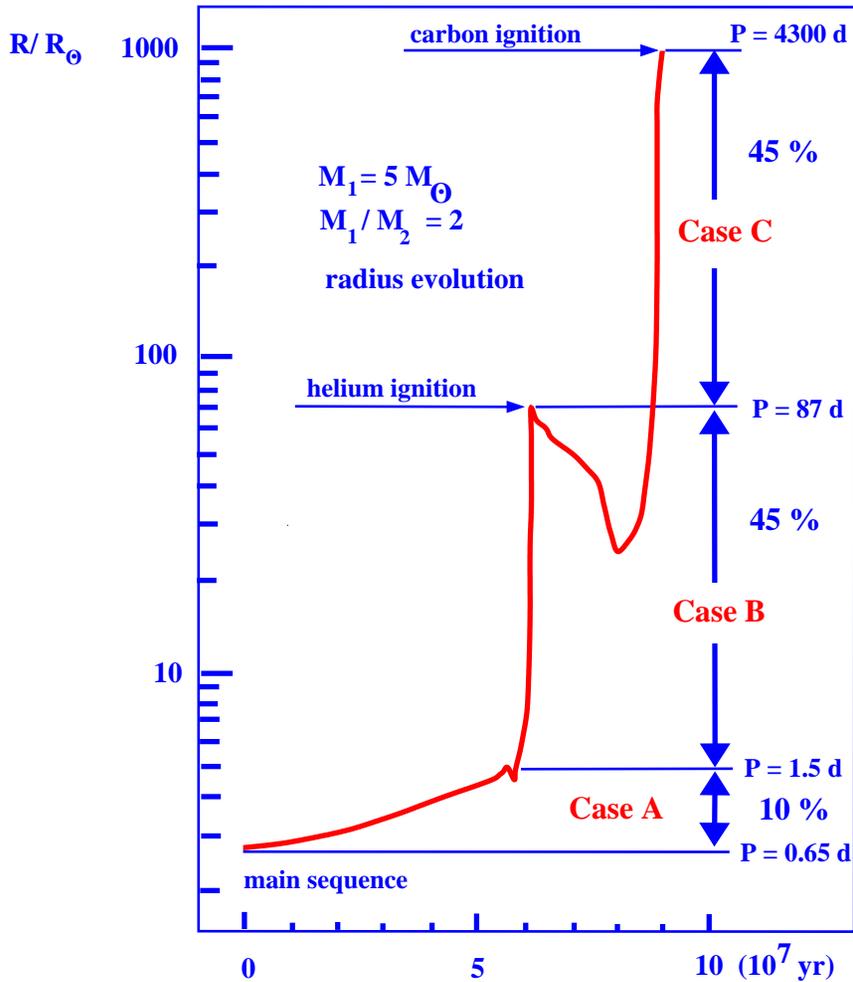
▷ if *envelope ejected* → *very close binary* (compact core + secondary)

▷ otherwise: *complete merger* of the binary components → *formation of a single, rapidly rotating star*

7.5 INTERACTING BINARIES (SELECTION) (Supplementary)

Classification of Roche-lobe overflow phases

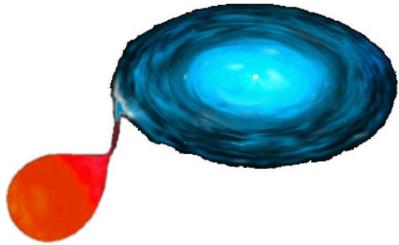
(Paczynski)



Algols and the Algol paradox

- *Algol* is an eclipsing binary with orbital period 69 hr, consisting of a B8 dwarf ($M = 3.7 M_{\odot}$) and a K0 subgiant ($M = 0.8 M_{\odot}$)
- the *eclipse* of the B0 star is *very deep* \rightarrow B8 star more *luminous* than the more evolved K0 subgiant
- the *less massive star is more evolved*
- inconsistent with stellar evolution \rightarrow *Algol paradox*
- explanation:
 - ▷ the K star was *initially the more massive star* and evolved more rapidly
 - ▷ *mass transfer* changed the mass ratio
 - ▷ because of the added mass the B stars becomes the more luminous component

Interacting binaries containing compact objects
(Supplementary)



- short orbital periods (11 min to typically 10s of days) → requires *common-envelope* and *spiral-in* phase

Cataclysmic Variables (CV)

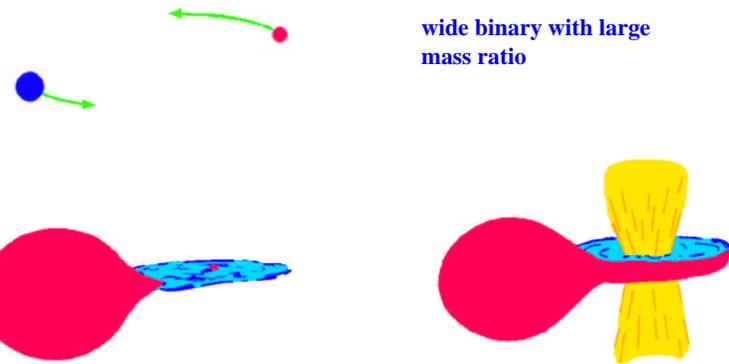
- main-sequence star (usually) transferring mass to a *white dwarf* through an *accretion disk*
- *nova outbursts: thermonuclear explosions* on the surface of the white dwarf
- orbit shrinks because of angular-momentum loss due to *gravitational radiation* and *magnetic braking*

X-Ray Binaries

- *compact component: neutron star, black hole*
- *mass donor* can be of low, intermediate or high mass
- very luminous *X-ray sources* (accretion luminosity)
- neutron-star systems: luminosity distribution peaked near the *Eddington limit*, (accretion luminosity for which radiation pressure balances gravity)

$$L_{\text{Edd}} = \frac{4\pi c G M}{\kappa} \simeq 2 \times 10^{31} \text{ W} \left(\frac{M}{1.4 M_{\odot}} \right)$$
- accretion of mass and angular momentum → *spin-up* of neutron star → formation of *millisecond pulsar*
- *soft X-ray transients: best black-hole candidates* (if $M_X > \text{max. neutron-star mass} \sim 2 - 3 M_{\odot}$ → likely black hole [but no proof of event horizon yet])

Formation of Low-Mass X-Ray Binaries (I)



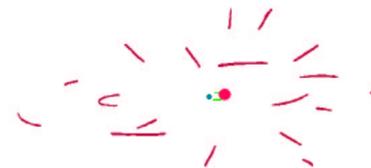
wide binary with large mass ratio



dynamical mass transfer



common-envelope and spiral-in phase



ejection of common envelope and subsequent supernova