BINARY STARS

- most stars are members of binary systems or multiple systems (triples, quadruples, quintuplets, ...)
- \bullet orbital period distribution: $P_{\rm orb} = 11\,min$ to $\sim 10^6\,yr$
- the majority of binaries are wide and do not interact strongly
- close binaries (with $P_{orb} \lesssim 10 \, yr$) can transfer mass \rightarrow changes structure and subsequent evolution
- approximate period distribution: $f(log P) \simeq const.$ (rule of thumb: 10% of systems in each decade of log P from 10^{-3} to $10^7 yr$)



• systems with eccentricities $\leq 10 d$ tend to be circular \rightarrow evidence for *tidal circularization*

Classification

- *visual binaries:* see the periodic wobbling of two stars in the sky (e.g. Sirius A and B); if the motion of only one star is seen: *astrometric binary*
- spectroscopic binaries: see the periodic Doppler shifts of spectral lines
 - \triangleright single-lined: only the Doppler shifts of one star detected
 - \triangleright *double-lined:* lines of both stars are detected
- *photometric binaries:* periodic variation of fluxes, colours, etc. are observed (caveat: such variations can also be caused by single variable stars: Cepheids, RR Lyrae variables)
- eclipsing binaries: one or both stars are eclipsed by the other one \rightarrow inclination of orbital plane $i \simeq 90^{\circ}$ (most useful for determining basic stellar parameters)

Radial Velocity (eccentric binaries)



for an eccentric binary

$$\begin{array}{l} \mathbf{x}(\mathbf{t}) = \mathbf{a} \; (\cos \mathbf{E} - \mathbf{e}) \\ \mathbf{y}(\mathbf{t}) = \mathbf{b} \; \sin \mathbf{E} \end{array} \qquad \qquad \\ \mathbf{tan} \frac{\mathbf{f}}{2} = \sqrt{\frac{1 + \mathbf{e}}{1 - \mathbf{e}}} \; \; \mathbf{tan} \frac{\mathbf{E}}{2} \end{array}$$

where the eccentric anomaly E is defined by Kepler's equation E - e sin E = $\frac{2\pi}{P}$ t = M (mean anomaly)

THE BINARY MASS FUNCTION

• consider a spectroscopic binary

• measure the radial velocity curve along the line of sight from $\frac{v_r}{c} \simeq \frac{\Delta \lambda}{\lambda}$ (Doppler shift) $M_1 \underset{i}{CM} \xrightarrow{v_1} \sum_{a_1} \sum_{a_2} M_2 \longrightarrow P = \frac{2\pi}{\omega} = 2\pi \frac{a_1 \sin i}{v_1 \sin i} = 2\pi \frac{a_2 \sin i}{v_2 \sin i}$ $\Rightarrow P = \frac{2\pi}{\omega} = 2\pi \frac{a_1 \sin i}{v_1 \sin i} = 2\pi \frac{a_2 \sin i}{v_2 \sin i}$ $\Rightarrow gravitational force = centripetal force$ $\Rightarrow \frac{GM_1M_2}{(a_1 + a_2)^2} = \frac{(v_1 \sin i)^2}{a_1 \sin^2 i} M_1, \quad \frac{GM_1M_2}{(a_1 + a_2)^2} = \frac{(v_2 \sin i)^2}{a_2 \sin^2 i} M_2$ substituting $(a_1 + a_2)^2 = a_1^2 (M_1 + M_2)^2 / M_2^2$, etc. $\Rightarrow f_1(M_2) = \frac{M_2^2 \sin^3 i}{(M_1 + M_2)^2} = \frac{P(v_1 \sin i)^3}{2\pi G}$ $f_2(M_1) = \frac{M_1^3 \sin^3 i}{(M_1 + M_2)^2} = \frac{P(v_2 \sin i)^3}{2\pi G}$

- f_1 , f_2 mass functions: relate observables $v_1 \sin i$, $v_2 \sin i$, P to quantities of interest M_1 , M_2 , $\sin i$
- measurement of f_1 and f_2 (for double-lined spectroscopic binaries only) determines $M_1 \sin^3 i$, $M_2 \sin^3 i$
 - \triangleright if i is known (e.g. for visual binaries or eclipsing binaries) \rightarrow $M_1,~M_2$
 - $\label{eq:main_star} \begin{array}{rcl} \triangleright \mbox{ for } & M_1 \ll M_2 & \rightarrow & f_1(M_2) \simeq M_2 \sin^3 i & (measuring $v_1 \sin i$ for star 1 constrains M_2) \end{array}$
- for *eclipsing binaries* one can also determine the *radii* of both stars (main source of accurate masses and radii of stars [and luminosities if distances are known])

The Roche Potential





THE ROCHE POTENTIAL

- restricted three-body problem: determine the motion of a test particle in the field of two masses M_1 and M_2 in a circular orbit about each other
- equation of motion of the particle in a frame rotating with the binary $\Omega = 2\pi/P$:

$$rac{\mathrm{d}^2ec{\mathbf{r}}}{\mathrm{d} \mathrm{t}^2} = -ec{
abla} \, \mathrm{U}_{\mathrm{eff}} - \underbrace{2ec{\mathbf{\Omega}} imes ec{\mathbf{v}}}_{\mathrm{Coriolis \ force}}$$

where the *effective potential* U_{eff} is given by

$$\mathbf{U}_{\mathrm{eff}} = -\frac{\mathbf{G}\mathbf{M}_{1}}{\left|\vec{r}-\vec{r}_{1}\right|} - \frac{\mathbf{G}\mathbf{M}_{2}}{\left|\vec{r}-\vec{r_{2}}\right|} - \underbrace{\frac{1}{2}\Omega^{2}\left(\mathbf{x}^{2}+\mathbf{y}^{2}\right)}_{\mathrm{centrifugal \ term}}$$

• Lagrangian points: five stationary points of the Roche potential U_{eff} (i.e. where effective gravity $\vec{\nabla}U_{eff} = 0$)

 \triangleright 3 saddle points: L₁, L₂, L₃

- Roche lobe: equipotential surface passing through the inner Lagrangian point L_1 ('connects' the gravitational fields of the two stars)
- approximate formula for the *effective Roche-lobe* radius (of star 2):

$$\mathbf{R_L} = \frac{0.49}{0.6 + \mathbf{q}^{2/3} \ln{(1 + \mathbf{q}^{-1/3})}} \mathbf{A}$$

where $\mathbf{q}=\mathbf{M}_1/\mathbf{M}_2$ is the mass ratio, A orbital separation.

Classification of close binaries

- Detached binaries:
 - both stars underfill their Roche lobes, i.e. the photospheres of both stars lie beneath their respective Roche lobes
 - gravitational interactions only
 (e.g. tidal interaction, see Earth-Moon system)
- Semidetached binaries:
 - \triangleright one star fills its Roche lobe
 - b the Roche-lobe filling component transfers matter to the detached component
 - \triangleright mass-transferring binaries
- Contact binaries:
 - > both stars fill or overfill their Roche lobes
 - \triangleright formation of a common photosphere surrounding both components
 - ▷ W Ursae Majoris stars

BINARY MASS TRANSFER

- 30 50% of all stars experience mass transfer by *Roche-lobe* overflow during their lifetimes (generally in late evolutionary phases)
- a) (quasi-)conservative mass transfer



 \triangleright mass loss + mass accretion

- $\triangleright \text{ the mass loser tends to} \\ \text{lose most of its} \\ \text{envelope} \rightarrow \text{formation} \\ \text{of } helium \ stars \\ \end{cases}$
- ▷ the accretor tends to be *rejuvenated* (i.e. behaves like a more massive star with the evolutionary clock reset)
- \triangleright orbit generally widens
- b) dynamical mass transfer \rightarrow common-envelope and spiral-in phase (mass loser is usually a red giant)



- > accreting component also fills its Roche lobe
- b mass donor (primary) engulfs secondary
- *spiral-in* of the core of the primary and the secondary immersed in a *common envelope*
- ightarrow if envelope ejected \rightarrow very close binary (compact core + secondary)
- \triangleright otherwise: complete merger of the binary components \rightarrow formation of a single, rapidly rotating star

Classification of Roche-lobe overflow phases



INTERACTING BINARIES (SELECTION)

Algols and the Algol paradox

- Algol is an eclipsing binary with orbital period 69 hr, consisting of a B8 dwarf $(M = 3.7 M_{\odot})$ and a K0 subgiant $(M = 0.8 M_{\odot})$
- the eclipse of the B0 star is very deep $\rightarrow B8$ star more luminous than the more evolved K0 subgiant
- the less massive star is more evolved
- inconsistent with stellar evolution \rightarrow *Algol paradox*
- explanation:
 - b the K star was initially the more massive star and evolved more rapidly
 - \triangleright mass transfer changed the mass ratio
 - \triangleright because of the added mass the B stars becomes the more luminous component

Interacting binaries containing compact objects

Formation of Low-Mass X-Ray Binaries (I)



ejection of common envelope and subsequent supernova



• short orbital periods (11 min to typically 10s of days) \rightarrow requires *common-envelope* and *spiral-in* phase

Cataclysmic Variables (CV)

- main-sequence star (usually) transferring mass to a *white dwarf* through an *accretion disk*
- nova outbursts: thermonuclear explosions on the surface of the white dwarf
- dwarf novae: accretion-disk instabilities
- orbit shrinks because of angular-momentum loss due to gravitational radiation and magnetic braking

X-Ray Binaries

- compact component: neutron star, black hole
- mass donor can be of low, intermediate or high mass
- very luminous *X-ray sources* (accretion luminosity)
- neutron-star systems: luminosity distribution peaked near the *Eddington limit*, (accretion luminosity for which radiation pressure balances gravity) $L_{Edd} = \frac{4\pi cGM}{\kappa} \simeq 2 \times 10^{31} \, W \, \left(\frac{M}{1.4 \, M_\odot}\right)$
- accretion of mass and angular momentum \rightarrow spin-up of neutron star \rightarrow formation of millisecond pulsar
- soft X-ray transients: best black-hole candidates (if $M_X > max$. neutron-star mass $\sim 2-3 M_{\odot} \rightarrow likely$ black hole [but no proof of event horizon yet])